Phase structure and thermodynamics of strongly-interacting matter

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in collaboration with

Pedro Costa, Eduardo S. Fraga, Lisa M. Haas, Hubert Hansen, Tina K. Herbst, Bruno W. Mintz, Mario Mitter, Jan M. Pawlowski, Rudnei O. Ramos, Jürgen Schaffner-Bielich & Andreas Zacchi

CPHT, École polytechnique; 28/09/2017









Introduction



2 Theoretical framework

3 Results

4 HIC phenomenology

6 Challenges



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Phase structure of the strong interaction

$$\mathcal{L}_{\text{QCD}} = \overline{\boldsymbol{q}} \left[i \gamma_{\mu} \left(\partial^{\mu} - i g A^{\mu} \right) - m + \gamma_0 \mu_f \right] \boldsymbol{q} - \frac{1}{4} G^a_{\mu\nu} G^{\mu\nu}_a$$



GSI Helmholtzzentrum für Schwerionenforschung

- low temperature & density: phase of confined, massive constituent quarks
- high temperature and/or density: phase of 'massless', deconfined quarks and gluons

Introduction

HIC phenomenology

Challenges

Conclusions

Observations on the Phase Diagram of QCD



Probed in

- Early universe at small density and high temperature
- · Compact star matter at small temperature and high density
- Relativistic heavy-ion collisions at LHC, RHIC, NICA, FAIR, ...

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Theoretical Insights to the Phase Diagram of QCD

$$\mathcal{L}_{\text{QCD}} = \overline{\boldsymbol{q}} \left[i \gamma_{\mu} \left(\partial^{\mu} - i g A^{\mu} \right) - m + \gamma_0 \mu_f \right] \boldsymbol{q} - \frac{1}{4} G^a_{\mu\nu} G^{\mu\nu}_a$$



→ Effective model capturing the major properties ~ chiral and centre symmetry

Phase diagram of the Polyakov-Quark-Meson model



RS and J. Schaffner-Bielich, Phys. Rev. D 93, 094014, 2016

→ There is a (small) region of a first order phase transition at large chemical potentials

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Phase diagram of the Polyakov-Quark-Meson model



RS, E. S. Fraga and J. Schaffner-Bielich, arXiv:1307.2851v1 [hep-ph]

→ There is a (small) region of a first order phase transition at large chemical potentials . . . which shrinks with increasing isospin imbalance

Effective Model for QCD

Major properties of strongly interacting matter $\,\, \sim \,\,$ symmetries

- Dynamical mass-generation of constituent quarks: spontaneous chiral symmetry breaking
- Deconfinement: spontaneous centre symmetry breaking Order parameter: Polyakov loop Φ
 ↔ Free energy of a test quark: Φ ~ exp(-F_q/T)
- → low temperature & density: chiral symmetry broken, centre symmetric high temperature and/or density: chiral symmetry restored, centre symmetry broken

 \Rightarrow Polyakov-loop Quark-Meson model, PNJL model

Ingredients of the Polyakov–Quark-Meson model

$$\mathcal{L}_{PQM} = \bar{q} \left[i \gamma_{\mu} \left(\partial^{\mu} - i A^{\mu} \delta_{\mu 0} + \mu_{f} \delta^{\mu 0} \right) - g \frac{\lambda_{a}}{2} \left(\sigma_{a} + i \gamma_{5} \pi_{a} \right) \right] q \\ + \frac{1}{2} \left(\partial_{\mu} \sigma_{a} \partial^{\mu} \sigma_{a} + \partial_{\mu} \pi_{a} \partial^{\mu} \pi_{a} \right) - U \left(\sigma_{a}, \pi_{a} \right) - \mathcal{U} \left(\Phi \left[A_{0} \right], \bar{\Phi} \left[A_{0} \right]; T \right)$$

Ingredients:

- constituent quarks
- scalar and pseudoscalar mesons $\langle ar{q}q
 angle$
 - → generation of constituent quark masses by meson exchange (Yukawa coupling): $m_f = g\sigma_f$
 - + vector mesons
- gauge fields \leftrightarrow Polyakov loop: $\Phi \sim \exp\left(i\int_0^\beta d\tau A_0\right)$ \rightarrow confinement (of quarks)
- → scalar fields: order parameters for chiral symmetry breaking
 → Polyakov-loop: order parameter for confinement

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Constraining the Phase diagram of the PQM model



RS, E. S. Fraga and J. Schaffner-Bielich, arXiv:1307.2851v1 [hep-ph]

HIC phenomenology

Constraints from the lattice @ $\mu = 0$

Order parameters



Centre symmetry

Challenges



L. Haas, RS, J. Braun, J. Pawlowski and J. Schaffner-Bielich, Phys. Rev. D 87, 076004, 2013 T. K. Herbst, M. Mitter, J. Pawlowski, B.-J. Schaefer and RS. Phys. Lett. B 731, 248-256, 2014

- Quantitative agreement in the chiral sector
- Polyakov loop of effective model $\Phi[\langle A_0 \rangle]$ upper limit of lattice $\langle \Phi \rangle$:

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 $\begin{array}{l} \Phi\left[\left\langle A_{0}\right\rangle\right]\geq\left\langle\Phi\left[A_{0}\right]\right\rangle\\ \textit{J. Braun, H. Gies and J. Pawlowski, Phys. Lett. B 684, 262-267, 2010\\ Phase structure and thermodynamics of strongly-interacting matter \end{array}$ 9

HIC phenomenology

Challenges

Constraints from the lattice @ $\mu = 0$



L. Haas, RS, J. Braun, J. Pawlowski and J. Schaffner-Bielich, Phys. Rev. D 87, 076004, 2013 T. K. Herbst, M. Mitter, J. Pawlowski, B.-J. Schaefer and RS. Phys. Lett. B 731, 248-256, 2014

Quantitative agreement or at least within the trend of lattice data

Seems to capture the important physics

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Constraining the Phase diagram of the PQM model



RS, E. S. Fraga and J. Schaffner-Bielich, arXiv:1307.2851v1 [hep-ph]

Results

HIC phenomenology

Challenges Conclusions

Constraints from the lattice $(Q) \mu \lesssim T_c$



RS, J. Schaffner-Bielich, Phys. Rev. D 93, 094014 (2016)

Curvature of the phase transition line: $T_{\rm c}(\mu) \propto \mu^2$ for $\mu \lesssim T_{\rm c}$

Comparison to chemical freeze-out of hadrons in relativistic HICs

\Rightarrow Within the lattice uncertainty, and consistent with experimental results (?)

 \rightarrow Impact on curvature: coupling strength of vector mesons

Constraining the Phase diagram of the PQM model



RS, E. S. Fraga and J. Schaffner-Bielich, arXiv:1307.2851v1 [hep-ph]

Constraints from Compact Star Masses and Radii

Calculate the equation of state within the model and solve the GR equations to gain the mass-radius relation.



A. Zacchi, RS, J. Schaffner-Bielich; Phys. Rev. D 92 (2015) 4, 045022

- Precise mass measurements
- Currently large uncertainties on radius constraints, high precision with future x-ray and gravitational wave measurements

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Constraining the Phase diagram of the PQM model



RS, E. S. Fraga and J. Schaffner-Bielich, arXiv:1307.2851v1 [hep-ph]

Input to HIC phenomenology

· Medium dependence of quark and meson masses



R. Marty and J. Aichelin, Phys. Rev. C 87, 034912 (2013)

- \rightarrow elastic scattering and hadronisation cross-sections
- → input to transport calculation
 - · transverse momentum and rapidity distributions
 - · elliptic flow as function of transverse momentum and rapidity
- \rightarrow transport coefficients:

shear & bulk viscosity, thermal conductivity

A. Abhishek, H. Mishra, S. Ghosh, arXiv:1709.08013 [hep-ph]

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Medium dependence of quark and meson masses
 → elastic scattering and hadronisation cross-sections



R. Marty, E. Bratkovskaya, W. Cassing and J. Aichelin, Phys. Rev. C 92, 015201 (2015), R. Marty and J. Aichelin, Phys. Rev. C 87, 034912 (2013)

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Introduction Theoretical framework Results HIC phenomenology (Challenges) Conclusions **A closer look to the Polyakov-loop potential** $\mathcal{L}_{QCD} = \bar{q} \left[i \gamma_{\mu} \left(\partial^{\mu} - i g A^{\mu} \right) + \gamma_{0} \mu_{f} - m \right] q - \frac{1}{4} G^{a}_{\mu\nu} G^{\mu\nu}_{a}$

$$\mathcal{L}_{PQM} = \bar{q} \left[i \gamma_{\mu} \left(\partial^{\mu} - i A^{\mu} \delta_{\mu 0} + \mu_{f} \delta^{\mu 0} \right) - g \frac{\lambda_{a}}{2} \left(\sigma_{a} + i \gamma_{5} \pi_{a} \right) \right] q \\ + \frac{1}{2} \left(\partial_{\mu} \sigma_{a} \partial^{\mu} \sigma_{a} + \partial_{\mu} \pi_{a} \partial^{\mu} \pi_{a} \right) - U \left(\sigma_{a}, \pi_{a} \right) - \mathcal{U} \left(\Phi[A_{0}], \bar{\Phi}[A_{0}]; T \right) \\ \mathcal{U} \left(\Phi, \bar{\Phi}; T \right) \leftrightarrow G_{\mu\nu} G^{\mu\nu} : \qquad \text{gluon interaction}$$

Parametrised Polyakov-loop potential fitted to pure gauge / Yang-Mills simulations

 $\Rightarrow \mathcal{U}_{\text{YM}}$: pure gauge / Yang-Mills / quenched Polyakov-loop potential !

 \rightarrow How does it change in the presence of dynamical quarks?

Unquenching the Polyakov-loop potential

Yang-Mills and QCD glue potential in the Functional Renormalisation Group



 $\Rightarrow \mathcal{U}_{glue}(t, \Phi) = \mathcal{U}_{YM}(t_{YM}(t), \Phi), \quad t_{YM}(t_{glue}) \simeq 0.57 t_{glue}$

L. Haas, RS, J. Braun, J. Pawlowski and J. Schaffner-Bielich, Phys. Rev. D 87, 076004, 2013

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Impact of unquenching: $\mathcal{U}_{\text{YM}} \rightarrow \mathcal{U}_{\text{glue}}$



L. Haas, RS, J. Braun, J. Pawlowski and J. Schaffner-Bielich, Phys. Rev. D 87, 076004, 2013 T. K. Herbst, M. Mitter, J. Pawlowski, B.-J. Schaefer and RS, Phys. Lett. B 731, 248-256, 2014

→ smoothens the transition



Results with different parameterisations ... differ



Different parameterisations ... differ already in YM



L. Haas, RS, J. Braun, J. Pawlowski and J. Schaffner-Bielich, Phys. Rev. D 87, 076004, 2013 RS et al., in preparation



Input to construct consistent P.-loop potentials

Polyakov-loop potential in continuum calculations



 $V(T, r_3, 0)/T^4$



U. Reinosa, J. Serreau, M. Tissier and N. Wschebor, Phys. Rev. D 93 (2016) 105002



Compare parameterisations and calculations



→ Calculated potentials flatter at t < 0⇒ Effect in PQM/PNJL qualitatively as unquenching

 \rightarrow Small height of barrier at phase transition

Introduction Theoretical framework Results HIC phenomenology Challenges Conclusions

- Conclusions
 - PQM/PNJL: Effective model to describe chiral symmetry breaking and confinement aspects
 - · Constrain framework and test improvements
 - against lattice data at vanishing density, small density and nonzero isospin density
 - astrophysics observations at large density, zero temperature
 - Make predictions on the existence and location of a CEP
 - Can serve as input to HIC phenomenology
 - Gauge sector requires revision: use input of ab-initio continuum calculations



HIC phenomenology

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Conclusions

Thank You for your attention!

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