

A connection between SYK physics and Glassy Dynamics

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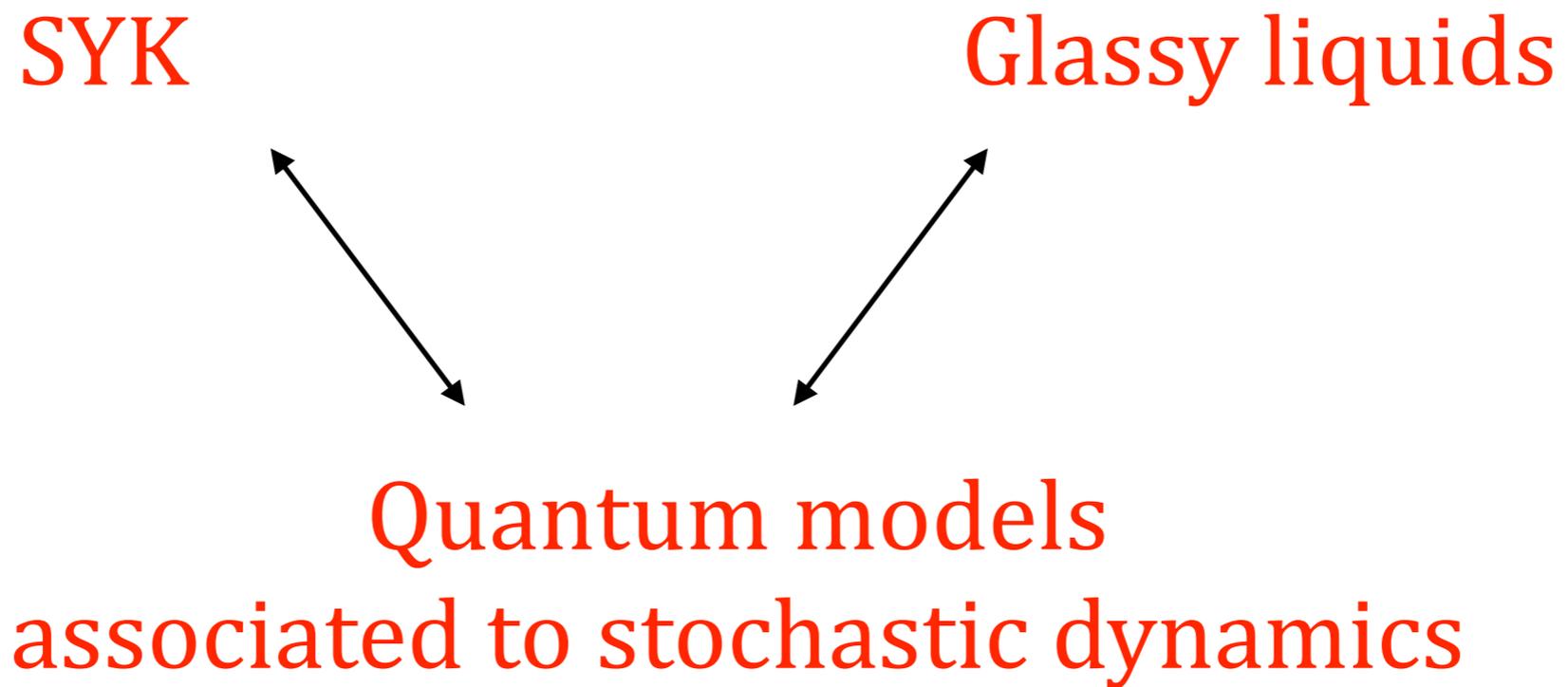
Davide Facoetti (ENS), Jorge Kurchan (ENS), David Reichman (Columbia)

Phys. Rev. B 2019, arXiv 1906.09228

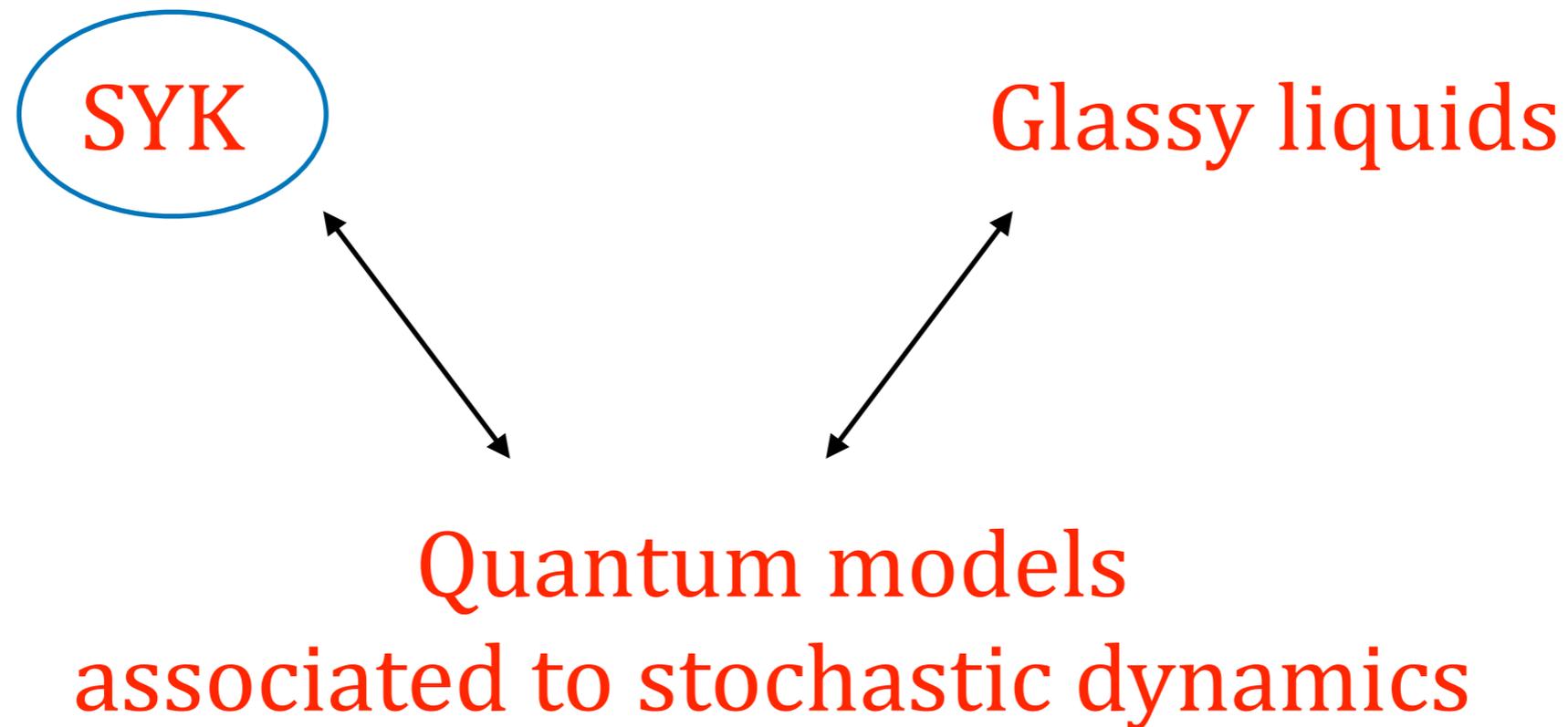


SIMONS FOUNDATION

A connection between glassy dynamics and SYK physics



A connection between glassy dynamics and SYK physics



SYK Model

Sachdev-Ye '93
Parcollet-Sachdev-Georges '01
Kitaev '15

$$H = \sum_{1 \leq i < j < k < l \leq N} J_{ijkl} \chi_i \chi_j \chi_k \chi_l$$

Majorana fermions $\{\chi_i, \chi_j\} = \delta_{ij}$

Quenched disorder

$$P(J_{ijkl}) \sim \exp(-N^3 J_{ijkl}^2 / 12J^2)$$

Studied in several contexts

- Related to models of extremal black holes
- Chaos bound: Maximally chaotic
- Non-Fermi liquid behavior

SYK main properties

Kitaev '15
Maldacena, Stanford '16
Polchinski, Rosenhaus '16

- Exactly solvable $G(\tau, 0) = \frac{1}{N} \sum_i \langle T \chi_i(\tau) \chi_i(0) \rangle$

$$\frac{1}{G(\omega)} = -i\omega - \Sigma(\omega), \quad \Sigma(t_1, t_2) = \text{Diagram: a circle with a horizontal line through its center. The top and bottom of the circle have small squares. The left and right ends of the horizontal line have small squares. The center of the circle is labeled 'k'. The top and bottom of the circle are labeled 'j'. The left and right ends of the horizontal line are labeled 'i'.$$

$$\frac{\partial G(t_1 - t_2)}{\partial t_1} + J \int_0^\beta dt G(t_1, t) G(t, t_2)^3 = -\delta(t_1 - t_2)$$

- Thermodynamics

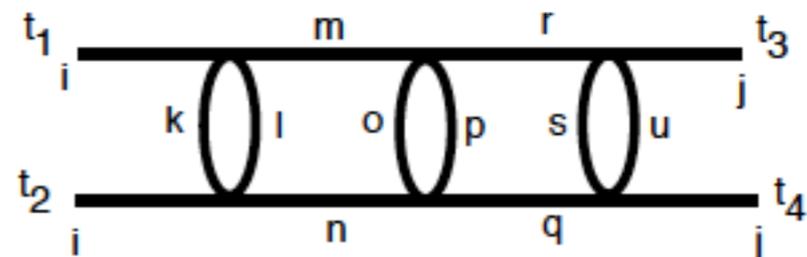
No ordering: Liquid
Finite entropy at zero temperature

- Critical at zero temperature

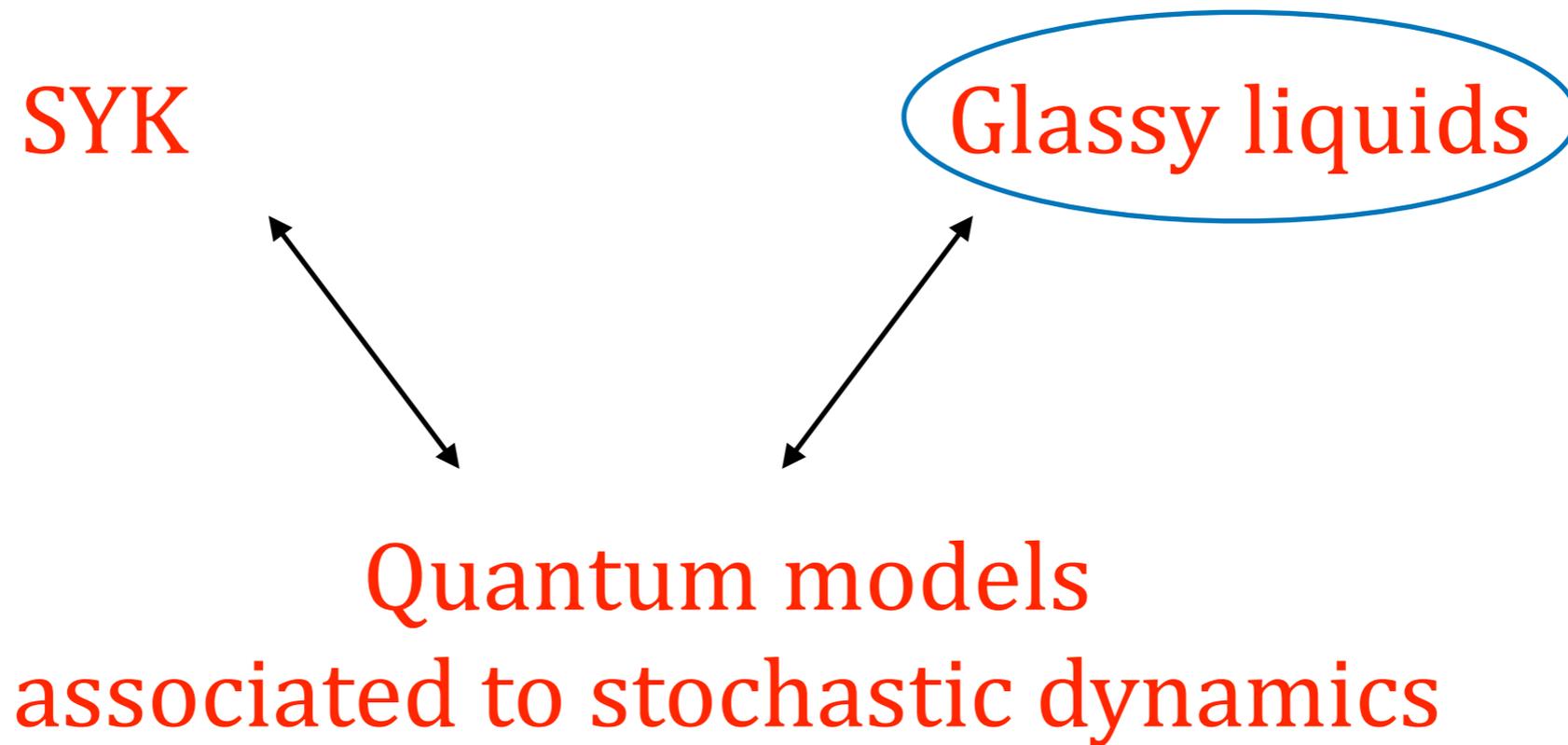
Gapless & power law behavior in time and in frequency
Time-scale set by $1/T$

- Low energy physics dominated by quasi time-reparametrization invariance

$$G(\tau, \tau') \rightarrow [f'(\tau) f'(\tau')]^{1/4} G(f(\tau), f(\tau'))$$



A connection between glassy dynamics and SYK physics



Exactly solvable models of classical glassy liquids

A key example

The p-spin spherical model

(p=4)

$$V = - \sum_{1 \leq i < j < k < l \leq N} J_{ijkl} S_i S_j S_k S_l$$

$$\sum_{i_1=1}^N S_{i_1}^2 = N$$

$$i_1 = 1, \dots, N$$

- Very Rough Energy Landscape

Exponential number of critical points

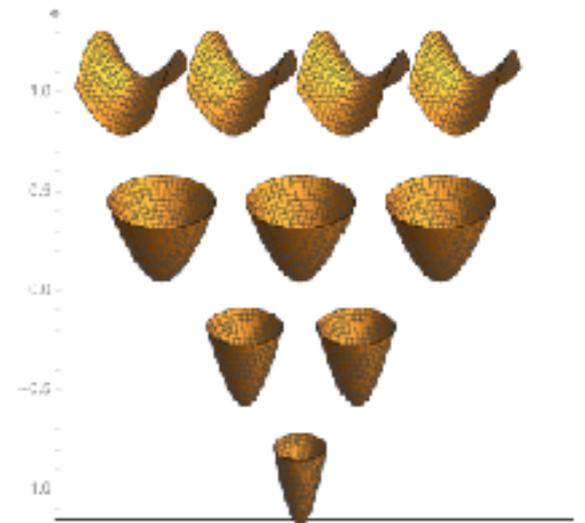
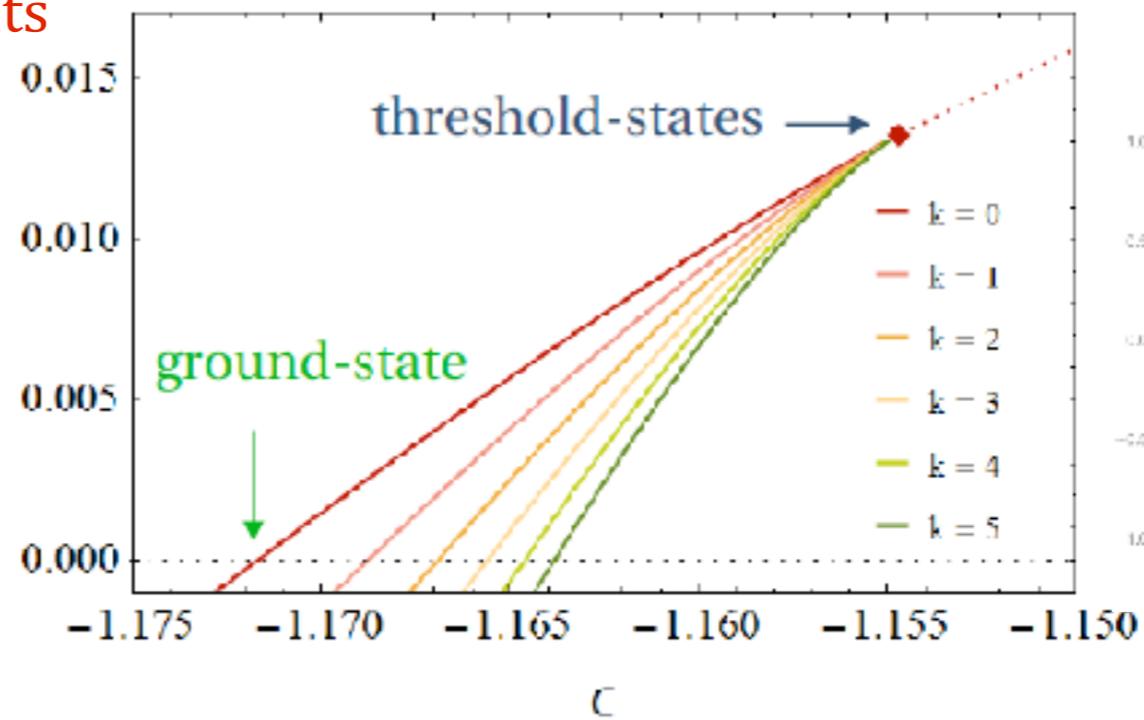
$$\mathcal{N}(e) \sim e^{N \Sigma_k(e)}$$

$\Sigma_k(e)$ Complexity of critical points of index k

Cavagna, Giardinà, Parisi (1998)

Auffinger, Ben Arous, Cerny (2013)

$$\Sigma_k(e) = \frac{1}{N} \ln \mathcal{N}(e)$$



Exactly solvable models of classical glassy liquids

- Their stochastic dynamics is exactly solvable by close eqs. on the two point function

$$\dot{S}(t) = -\nabla_S V + \eta(t) \quad \langle \eta_\alpha(t) \eta_\beta(t') \rangle = 2T \delta(t - t') \delta_{\alpha, \beta}$$

$$C(t, t') = \frac{1}{N} \sum_i \langle S_i(t) S_i(t') \rangle$$

Equilibrium dynamics $\dot{C}(\tau) = -TC(\tau) - \frac{P}{2T} \int_0^\tau du C^{p-1}(\tau - u) \dot{C}(u)$

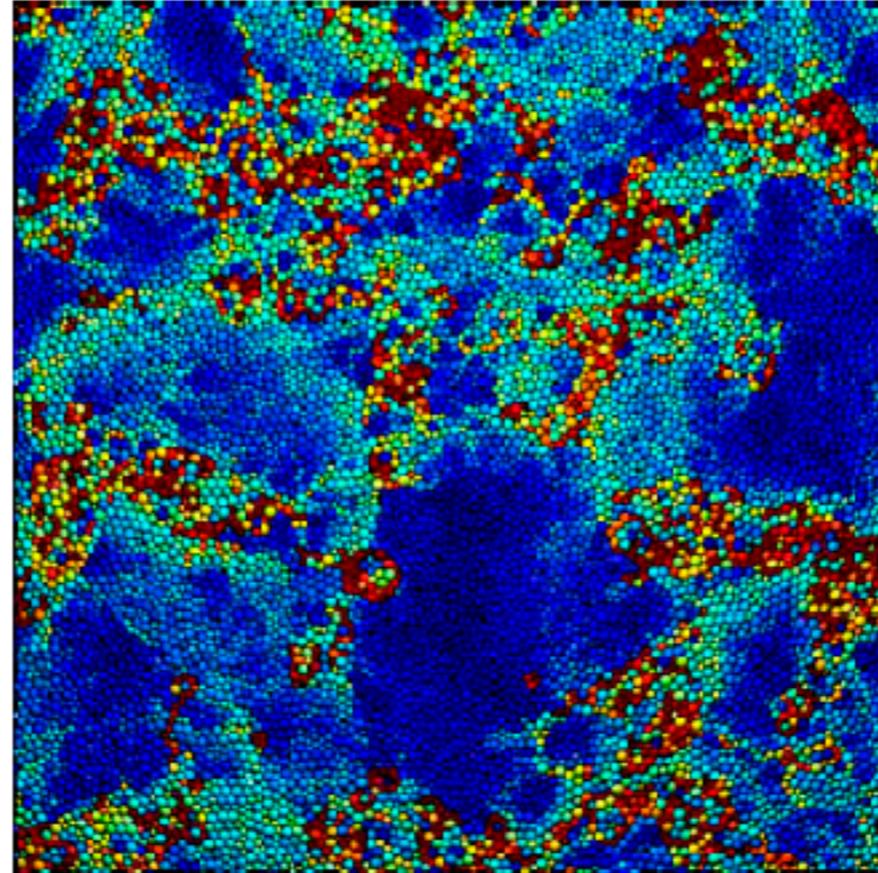
- At T_d dynamical glass transition; quenches below T_d aging
Gapless relaxation & power law behavior in time

Criticality of glassy dynamics

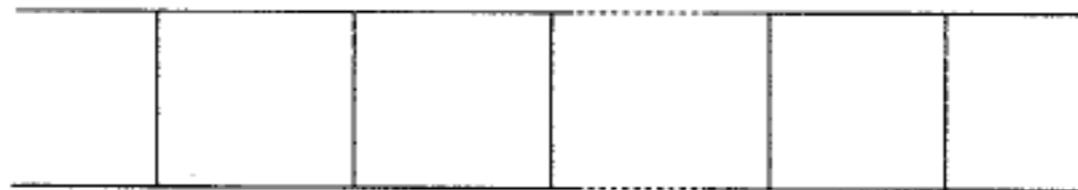
- Dynamical heterogeneity

The critical fluctuations of glassy dynamics

Biroli, Berthier, Bouchaud, Cipelletti, Van Saarloos
Oxford Univ Press 2011



Encoded in ladder diagrams:
dominant contribution to
4-point functions



Biroli, Bouchaud '04

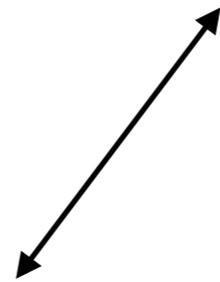
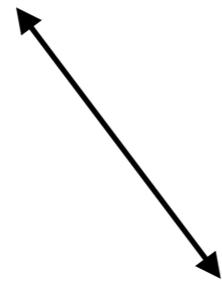
Related to quasi time-reparametrization invariance

Chamon, Cugliandolo et al '07

A connection between glassy dynamics and SYK physics

SYK

Glassy liquids



Quantum models
associated to stochastic dynamics

Mapping between stochastic and quantum dynamics

HE: stochastic quantisation

CMT: Rokhsar-Kivelson

SP: quantum-stochastic mapping

Langevin equation $\dot{q}_i(t) = -\frac{\partial V}{\partial q_i} + \eta_i(t)$

$$\partial_t P_t(\mathbf{q}) = \sum_i \frac{\partial}{\partial q_i} \left[T_s \frac{\partial}{\partial q_i} + \frac{\partial V}{\partial q_i} \right] P_t(\mathbf{q}) \equiv -H_{\text{FP}} P$$

Quantum Hamiltonian

$$H = \frac{T_s}{2} e^{V/2T_s} H_{\text{FP}} e^{-V/2T_s} = \sum_i \left[-\frac{T_s^2}{2} \frac{\partial^2}{\partial q_i^2} + \boxed{\frac{1}{8} \left(\frac{\partial V}{\partial q_i} \right)^2 - \frac{T_s}{4} \frac{\partial^2 V}{\partial q_i^2}} \right]$$

$\swarrow V_{\text{eff}}(\mathbf{q})$

- Zero ground-state energy & ground state \rightarrow Boltzmann-Gibbs distribution

$$\psi_{\text{GS}} = \frac{e^{-\frac{V}{2T_s}}}{\sqrt{Z}}$$

Mapping between stochastic and quantum dynamics

- Mapping between ground state quantum dynamics and stochastic dynamics

$$\frac{1}{N} \sum_i \langle S_i(\tau) S_i(0) \rangle_{clas} = \frac{1}{N} \sum_i \langle S_i(\tau) S_i(0) \rangle_{Qua} = \int_0^\infty \frac{d\omega}{2\pi} \rho(\omega) e^{-\omega\tau}$$

$$\frac{1}{2N} \sum_i \langle \{S_i(t), S_i(0)\} \rangle_{Qua,R} = \int_0^\infty \frac{d\omega}{2\pi} \rho(\omega) \cos \omega t$$

Henley '04

Biroli, Chamon, Zamponi '08

criticality in stochastic dynamics \leftrightarrow criticality in quantum dynamics

- Eigenvectors with eigenvalues much smaller than T_q \rightarrow metastable states
with lifetime much larger than β_q

Gaveau-Schulman '00

Bovier '00

Biroli, Kurchan '01

The quantum model

P-spin spherical model

$$V = - \sum_{1 \leq i < j < k < l \leq N} J_{ijkl} S_i S_j S_k S_l$$
$$\sum_{i_1=1}^N S_{i_1}^2 = N$$
$$i_1 = 1, \dots, N$$

Stochastic-Quantum Mapping



$$H = \sum_i -\frac{T_s^2}{2} \frac{\partial^2}{\partial S_i^2} + V_{eff}$$
$$V_{eff} = \sum_i \left[\frac{1}{8} \left(\frac{\partial V}{\partial S_i} \right)^2 - \frac{T_s}{4} \frac{\partial^2 V}{\partial S_i^2} \right]$$

Quantum model from a parent stochastic glassy dynamics

Exactly solvable by close eqs. on the two point functions

Results

Quantum Model

Parent stochastic glassy dynamics

- Thermodynamics

Finite entropy at zero temperature ←

Exponential number of metastable states with exponentially diverging life-time

$$\lim_{\beta_q \rightarrow \infty} \lim_{N \rightarrow \infty} \ln \text{Tr} e^{-\beta_q H} = S = \ln \mathcal{N} = N s_{meta}$$

No ordering: Quantum Liquid ←

Metastable states with finite life-time

Results

Quantum Model

Parent stochastic glassy dynamics

- Criticality at zero temperature

Gapless & power law behavior in time
cut-off by temperature

Dynamic criticality of glassy dynamics
Wide distribution of relaxation times



$$\rho_q(\omega) = \rho_c(1/\tau)$$

Results

Quantum Model

• Low energy physics dominated by quasi time-reparametrization invariance

Parent stochastic glassy dynamics

Quasi time-reparametrization invariance of glassy dynamics

$$t \rightarrow h(t)$$



Results

Quantum Model

- Thermodynamics

No ordering: Quantum Liquid
Finite entropy at zero temperature



Parent stochastic glassy dynamics

Exponential number of metastable states with exponentially diverging life-time

- Critical at zero temperature

Gapless & power law behavior in time cut-off by temperature



Dynamic criticality of glassy dynamics
Wide distribution of relaxation times

- Low energy physics dominated by quasi time-reparametrization invariance



Quasi time-reparametrization invariance of glassy dynamics

Differences: purely bosonic model, critical exponents, time-reparametrization transformations of correlation functions

Conclusion

Quantum models from parent stochastic glassy dynamics have SYK physics

Realised in many models: p-spin spherical model and many others (quantum spins)

Glassy models without disorder

Witten's '16 model for SYK without disorder

Transition from glassy to simple dynamics

Banerjee, Altman '17 model for transition to non-SYK physics

Open questions

Differences with SYK

Quantum chaotic properties

Effective theory of glassy dynamics & the Schwarzian action for the soft-modes