

Bulk Reconstruction Beyond the Entanglement Wedge

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(1911.00519)

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Outline

- ① Introduction: Subregion Duality in AdS/CFT
- ② Minimal Boundary-Anchored Surfaces
- ③ Bulk Metric Reconstruction
- ④ Discussion: Implications for Subregion Duality
- ⑤ Summary

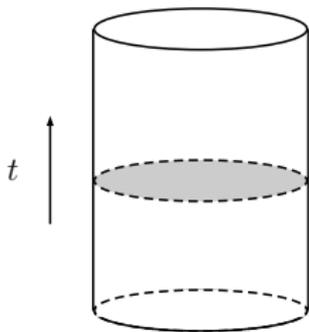
The AdS/CFT correspondence

(*certain*) states in (*certain*)
conformal field theories
(*in the right limit*)

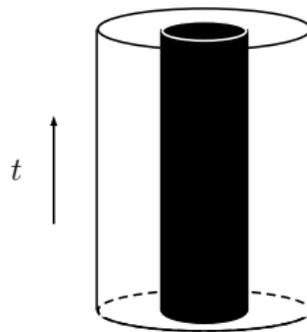
\leftrightarrow

asymptotically locally
Anti de Sitter geometry
(*aAdS*)

Examples:



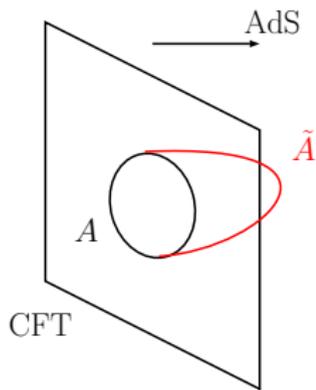
$|0\rangle$



$$\rho = \frac{1}{Z} \sum_n e^{-\beta E_n} |n\rangle \langle n|$$

The Ryu-Takayanagi formula

$$S(\rho_A) = \frac{\text{area}(\tilde{A})}{4G} + O(1)$$

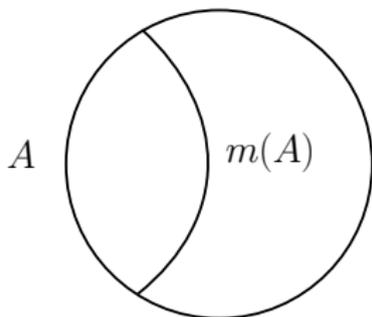


$\tilde{A} \equiv m(A) \equiv$ minimal surface homologous to A

Note: We'll stick to the static case throughout, except for briefly at the end.

The entanglement wedge

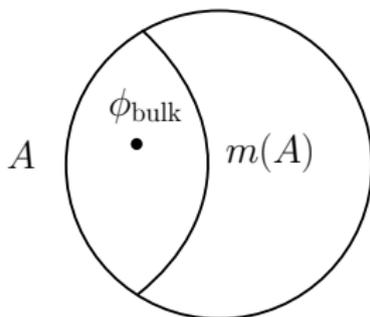
Definition: The *entanglement wedge* of A is the bulk region interior to $A \cup m(A)$.



→ subregion-subregion duality

Evidence for subregion duality*: HKLL

- Hamilton, Kabat, Lifschytz, Lowe (2006); hep-th/0606141
- Morrison (2014); 1403.3426
- others/follow-ups

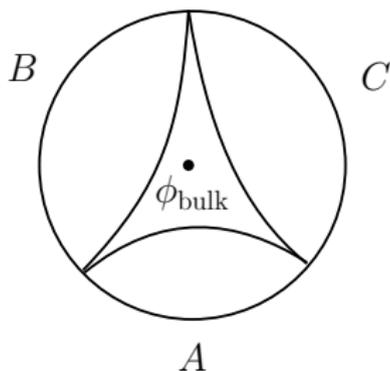


$$\phi_{\text{bulk}}(x, z) = \int_A dx' K(z, x|x') \mathcal{O}_{\text{boundary}}(x')$$

*for operators in the *causal wedge*

Evidence for subregion duality*: quantum error correction

- Almheiri, Dong, Harlow (2014); 1411.7041
- Dong, Harlow, Wall (2016); 1601.05416



c.f. qutrit code

$$|\tilde{0}\rangle = \frac{1}{\sqrt{3}} (|000\rangle + |111\rangle + |222\rangle)$$

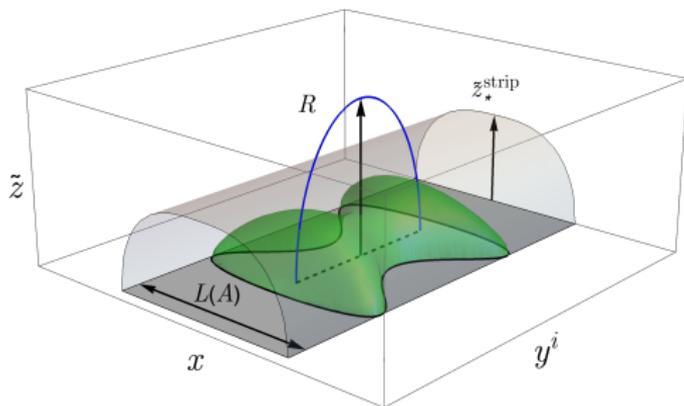
$$|\tilde{1}\rangle = \frac{1}{\sqrt{3}} (|012\rangle + |120\rangle + |201\rangle)$$

$$|\tilde{2}\rangle = \frac{1}{\sqrt{3}} (|021\rangle + |102\rangle + |210\rangle)$$

*for operators in the *entanglement wedge*

Tension with subregion duality

- ▶ In 3D gravity, there is only one type of minimal surface
- ▶ In higher dimensions, structure is richer



Observation: different dim minimal surfaces probe different bulk regions
→ let's exploit this observation

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Please forget about holography now.

Bulk reach of boundary-anchored minimal surfaces

The setup:

- ▶ static $d + 1$ dim'l asymptotically AdS space-time
- ▶ Poincaré patch,

$$ds^2 = \frac{1}{\tilde{z}^2} (-f(\tilde{z})dt^2 + dx_i dx^i + h(\tilde{z})d\tilde{z}^2)$$

- ▶ Boundary at $\tilde{z} = 0$, $f(\tilde{z}), g(\tilde{z}) \rightarrow 1$ as $\tilde{z} \rightarrow 0$
- ▶ Look at single space-like slice

B_k : simply-connected, k -dim'l submanifold in the boundary of the slice,
 $1 \leq k \leq d - 1$, $\partial B_k \neq \emptyset$

$m(B_k)$: minimal k -surface with $\partial m(B_k) = \partial B_k$

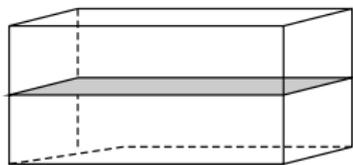
Q: deepest bulk reach \tilde{z}_* of $m(B_k)$?

Example: pure AdS, strip

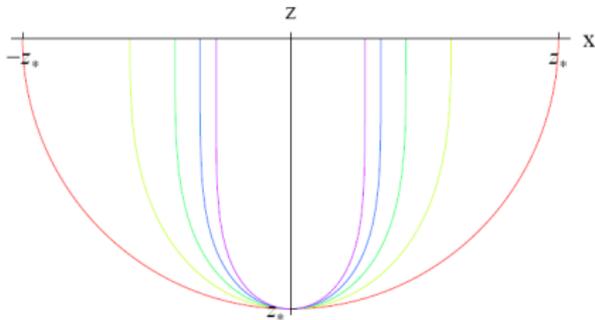
- Hubeny (2012); 1203.1044

$x \in [-L/2, L/2]$ and $y^i \in \mathbb{R}$, $i = 2, \dots, k$

$$\tilde{z}_*^{\text{strip}} = L \frac{(k-1) \Gamma(\frac{2k-1}{2k-2})}{\sqrt{\pi} \Gamma(\frac{k}{2k-2})}$$



e.g. 2d strip, 3d boundary



Example: pure AdS, ball

- Hubeny (2012); 1203.1044

$$\sum_{i=1}^k (x^i)^2 \leq R$$

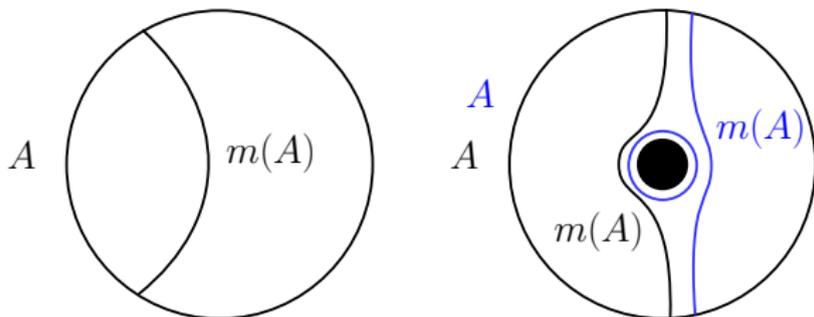
$$\tilde{z}_*^{\text{ball}} = R$$

k -wedges

Definition:

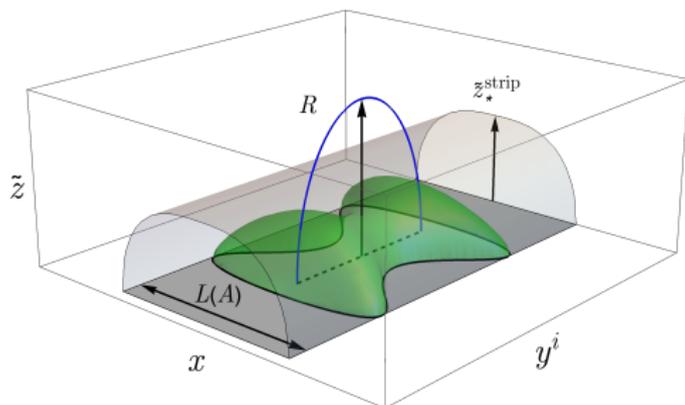
- ▶ A : simply-connected, $(d - 1)$ -dimensional subregion in boundary of static aAdS_{d+1} slice
- ▶ k -wedge of A , $W_k(A)$: set of all points that lie on a k -dimensional minimal surface $m(B_k)$ for at least one simply-connected submanifold $B_k \subseteq A$, where $1 \leq k \leq d - 1$.

Eg: $W_{d-1}(A) \equiv$ entanglement wedge of A (if no entanglement shadows.)



k -wedges reach beyond the entanglement wedge

- ▶ Ok, not too surprising!
- ▶ But, can construct situations where k -wedges probe parametrically beyond the entanglement wedge



- ▶ $d - 2$ directions y^i
- ▶ shown: “peanut”, $(d - 1)$ -dim'l strip, $(d - 2)$ -dim'l disk

k -wedges reach beyond the entanglement wedge

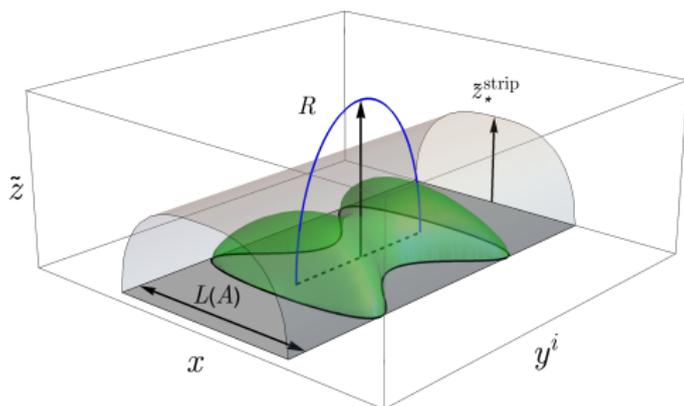
Exploit entanglement wedge nesting

- Akers, Koeller, Leichenauer, Levine (2016); 1610.08968

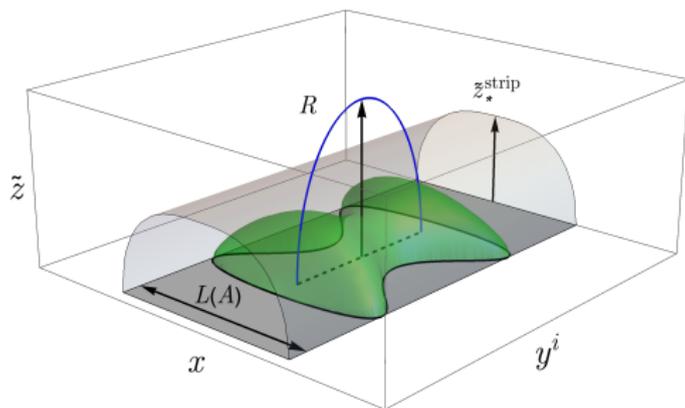
Theorem: If $A \subset A'$, then $W_E(A) \subset W_E(A')$.

$$\Rightarrow \tilde{z}^{\text{peanut}} \leq z^{\text{strip}}$$

- ▶ by choosing A appropriately, can inscribe k -balls with arbitrarily large R



Notes



- ▶ Can always push $W_{d-1}(A)$ arbitrarily close to boundary
- ▶ Subleading corrections known if you really want (Hubeny)

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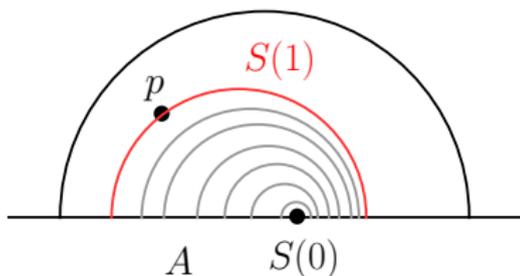
The BCFK Theorem

- Bao, Cao, Fischetti, Keeler (2019); 1904.04834

In a nutshell: $W_2(A)$ is metric-reconstructible

Theorem:

- ▶ Let $d \geq 3$
- ▶ A : $(d - 1)$ -dim'l topological ball
- ▶ Suppose for each $p \in W_2(A)$, \exists continuous family $m(S(\lambda))$
 - $\lambda \in [0, 1]$
 - $S(\lambda) \subset A$, topological disks
 - $S(0) \equiv$ single point in A
 - $S(1) = B : p \in m(B)$



Then, knowledge of the areas of all minimal 2-surfaces whose anchors lie in A guarantees a unique reconstruction of the bulk metric in $W_2(A)$.

Please remember holography now.

2-surface areas from boundary operators

Example: Wilson loops

- Maldacena (1998); *hep-th/9803002*
- Drukker, Gross, Ooguri (1999); *hep-th/9904191*
- ▶ Specify loop as boundary of some disk $B \subset A$

$$\langle W(\partial B) \rangle_{\rho_A} = \int_{\sigma \sim B} \mathcal{D}\sigma e^{-\sqrt{\lambda} S[\sigma]}$$

Large- N , large- λ :

- ▶ $S[\sigma]$ evaluates to (Legendre transform of) area of σ
- ▶ Saddle point approximation:

$$\langle W(\partial B) \rangle_{\rho_A} = e^{-\sqrt{\lambda} \mathcal{A}[m(B)]}$$

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Recap

- ▶ In $d \geq 4$, can construct boundary subregions A where $W_2(A)$ extends parametrically far beyond $W_{d-1}(A)$ (*entanglement wedge*)
- ▶ If there is *any* boundary operator localized to A which lets you learn the areas of minimal 2-surfaces $m(B) \subset W_2(A)$ (e.g. *Wilson loop expectation values*), then BCFK \Rightarrow learn the bulk metric in $W_2(A)$

The Alternatives:

- 1 In general, no boundary ops. in A calculated by minimal 2-surface areas.
- 2 ρ_A knows about more than just the entanglement wedge.

Option 1: No boundary operators reveal 2-surface areas

Or, “No boundary operators *localized to A* reveal *sufficiently large 2-areas*”

E.g. Wilson loop expectation values

- ▶ Cases where $\langle W \rangle$ does reveal minimal 2-surface areas, like pure AdS, must be “accidents”
- ▶ c.f. analyticity
 - know $g_{\mu\nu}$ in any open neighbourhood of any point, know $g_{\mu\nu}$ everywhere
- ▶ Either $\langle W \rangle$ breaks down at the mildest non-analyticity
 - e.g. gravitational wave with discontinuous 3rd derivative
- ▶ Or $\langle W \rangle$ fails for sufficiently large loops

Note: applies to any operator that reveals minimal 2-areas

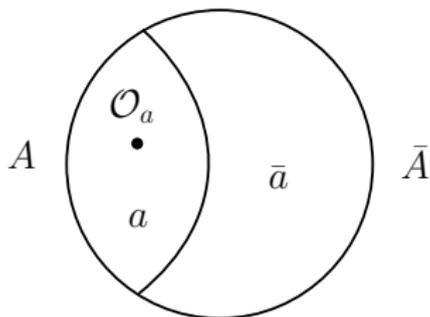
Option 2: Breakdown of subregion duality

Or, ρ_A knows about more than just the entanglement wedge.

Q: How is it still consistent with the error correction picture?
(It's a theorem, after all...)

Recall Dong/Harlow/Wall's "AdS/CFT as Quantum Error Correction"

$$\begin{aligned}\mathcal{H}_{\text{code}} &\subset \mathcal{H}_{\text{CFT}} \\ \mathcal{H}_{\text{CFT}} &= \mathcal{H}_A \otimes \mathcal{H}_{\bar{A}} \\ \mathcal{H}_{\text{code}} &= \mathcal{H}_a \otimes \mathcal{H}_{\bar{a}}\end{aligned}$$



Entanglement Wedge Reconstruction:

“any \mathcal{O}_a on \mathcal{H}_a represented by \mathcal{O}_A on \mathcal{H}_A ”

AdS/CFT as QEC

Theorem (Dong, Harlow, Wall)

- ▶ Let \mathcal{H} : $\dim \mathcal{H} < \infty$, $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_{\bar{A}}$, $\mathcal{H}_{\text{code}} \subset \mathcal{H}$
- ▶ Suppose $\mathcal{O}, \mathcal{O}^\dagger : \mathcal{H}_{\text{code}} \rightarrow \mathcal{H}_{\text{code}}$

If: $\forall |\phi\rangle, |\psi\rangle \in \mathcal{H}_{\text{code}}$, $\exists \mathcal{H}_{\text{code}} = \mathcal{H}_a \otimes \mathcal{H}_{\bar{a}}$ such that $\mathcal{O} = \mathcal{O}_a \otimes I_{\bar{a}}$

And If:

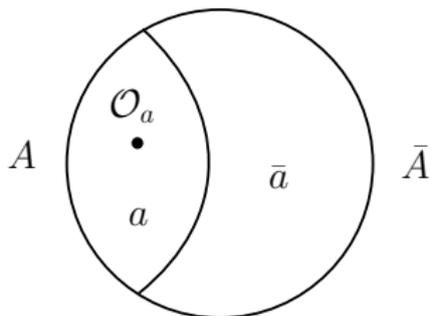
$$\begin{aligned}\rho_{\bar{A}} &= \text{Tr}_A |\phi\rangle\langle\phi| & \sigma_{\bar{A}} &= \text{Tr}_A |\psi\rangle\langle\psi| \\ \rho_{\bar{a}} &= \text{Tr}_a |\phi\rangle\langle\phi| & \sigma_{\bar{a}} &= \text{Tr}_a |\psi\rangle\langle\psi|\end{aligned}$$

satisfy $\rho_{\bar{a}} = \sigma_{\bar{a}} \Rightarrow \rho_{\bar{A}} = \sigma_{\bar{A}}$

Then:

- 1 For any $X_{\bar{A}} : \mathcal{H}_{\bar{A}} \rightarrow \mathcal{H}_{\bar{A}}$ and any $|\phi\rangle \in \mathcal{H}_{\text{code}}$, $\langle\phi|[\mathcal{O}, X_{\bar{A}}]|\phi\rangle = 0$
- 2 $\exists \mathcal{O}_A : \mathcal{H}_A \rightarrow \mathcal{H}_A$ such that $\mathcal{O}_A|\phi\rangle = \mathcal{O}|\phi\rangle \forall |\phi\rangle \in \mathcal{H}_{\text{code}}$

No conflict with Dong/Harlow/Wall

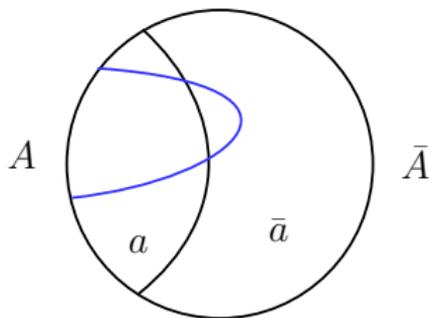


Subregion duality à la DHW \leftrightarrow sufficiently small algebra of bulk operators on a can be represented on A

An apparent puzzle regarding complementary recovery

For any $X_{\bar{A}} : \mathcal{H}_{\bar{A}} \rightarrow \mathcal{H}_{\bar{A}}$ and any $|\phi\rangle \in \mathcal{H}_{\text{code}}$, $\langle \phi | [\mathcal{O}, X_{\bar{A}}] | \phi \rangle = 0$

Untrue for, e.g., boundary Wilson loops?

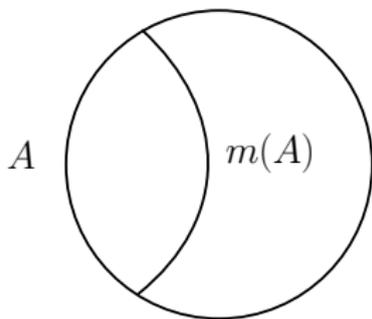


Two observations:

- 1 $\langle W \rangle \sim \text{area}$ not an operator equation
- 2 states with different metrics in \bar{a} not in same $\mathcal{H}_{\text{code}}$

Boundary reconstruction?

Given $W_E(A)$, what do you learn about ρ_A ?



Metric reconstruction from other data

- ▶ Can you reconstruct the metric from other minimal surfaces?
- ▶ What can you reconstruct in $W_k(A)$ for $k \neq 2$?

Example: Boundary-anchored geodesics

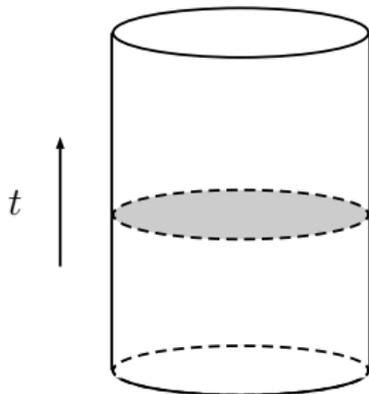
- ▶ Metric reconstruction known as *boundary rigidity problem* in applied maths community
 - Not proven for $A \subset \text{aAdS}$ boundary
- ▶ Holographically: calculated by 2-point functions in the geodesic approximation

$$\langle \mathcal{O}(x)\mathcal{O}(y) \rangle \sim e^{-|\text{geod}(x,y)|}$$

Covariant generalization?

Or, can you play the same games in non-static geometries?

E.g. BCFK still applies covariantly.



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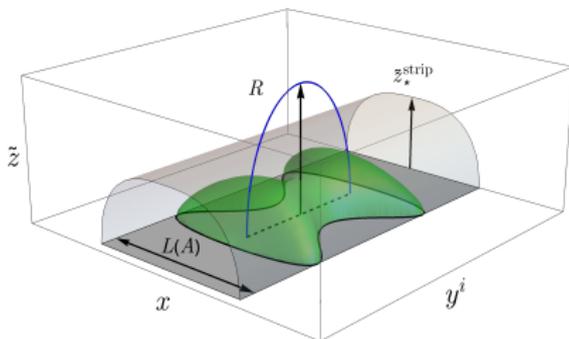
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Bulk reconstruction beyond the entanglement wedge

- 1 Given slice of static aAdS $_{d+1}$, can construct subregions A where $W_2(A)$ extends parametrically far beyond A 's entanglement wedge.
- 2 Spectrum of minimal 2-surfaces areas encodes bulk metric in $W_2(A)$
 - BCFK theorem
- 3 there are settings where minimal 2-surface areas are given by boundary operator $\langle \cdot \rangle$ localized to A
 - e.g. Wilson loops

\therefore can construct examples where ρ_A reveals bulk metric beyond its entanglement wedge

Alternatives and follow-up questions



- 1 No boundary operators in A reveal too large minimal 2-surface areas.
- 2 ρ_A knows about more than just the entanglement wedge.

So, which is it?

What does the entanglement wedge encode?

Other minimal surface data / non-static?