



UNIVERSITÀ  
DEGLI STUDI  
DI PADOVA

# Finding Axions in a Universe of Data and Envisioning Their Use as Multi-Messenger Probes

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the European Union

msca\_axitools

- Brief axion intro
- Axion global fits and model space
- The solar axion flux, its uncertainties, and the discovery potential for future helioscopes
- Post-discovery physics of axions

## The strong CP problem

The QCD Lagrangian contains the “ $\theta$  term”

$$\mathcal{L}_{\text{QCD}} = \dots - \frac{\alpha_s}{8\pi} \theta G_{\mu\nu}^a \tilde{G}^{\mu\nu,a} = \dots + \frac{\alpha_s}{2\pi} \theta \mathbf{E}^a \cdot \mathbf{B}^a$$

with gluon field dual  $\tilde{G}^{\mu\nu,a} \equiv \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} G_{\alpha\beta}^a$  and a constant  $\theta$ .

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- $\mathbf{E}^a \cdot \mathbf{B}^a \propto \partial_\alpha (\epsilon^{\alpha\beta\gamma\delta} A_\beta^a \partial_\gamma A_\delta^a)$  i.e. a total derivative, **but** also anomalous: can't be ignored due to instanton solutions
- $\alpha_s(m_Z) = 0.1183(9)$ <sup>2309.12986</sup>  $\rightarrow$  measure  $\theta$

## The strong CP problem

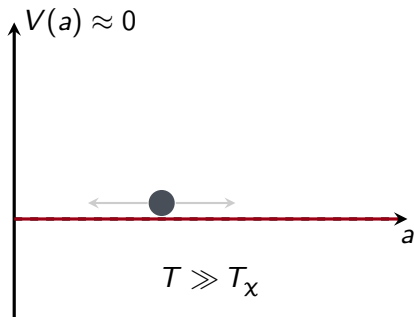
- $\theta \sim \mathcal{O}(1)$  should induce a measurable electric dipole moment of the neutron,  $d_n$
- Current bound:  $|d_n| < 1.8 \times 10^{-26} \text{ e cm (90\% CL)}$ <sup>2001.11966</sup>  
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- *Why is  $\theta$  so small?* Puzzling because ...
  - ...  $CP$  violation exists in weak interactions (kaon decay, 1964), EM doesn't have  $CP$ -violating diagrams (at tree level)
  - ... actually  $\theta \mapsto \theta - \arg \det(M_q)$ , so small  $\theta$  is even more surprising
  - ... all allowed terms should be  $\sim \mathcal{O}(1)$ !?

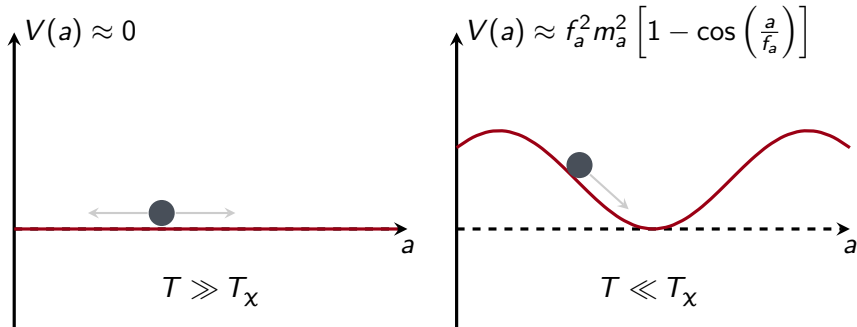
## Axion dark matter – realignment mechanism

- At early times,  $T \gg T_\chi \sim T_{\text{QCD,c}} = 158.1(5) \text{ MeV}$ ,<sup>2002.02821</sup> the axion field  $a$  can fluctuate freely



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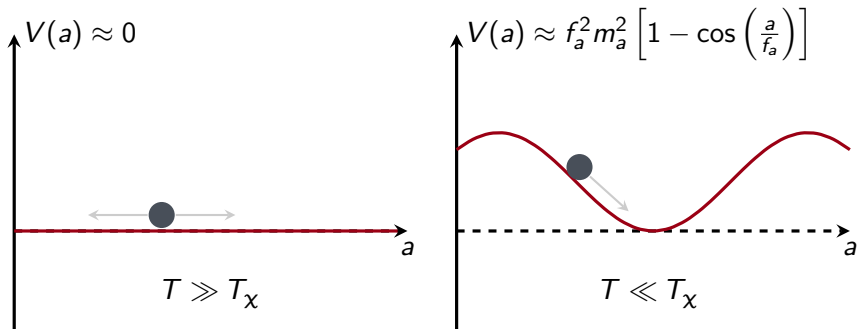
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- ➔ *Strong CP problem solved dynamically by promoting  $\theta \mapsto a/f_a$*
- ➔ *Oscillating scalar field behaves as DM*



Axion = pNGB from U(1) symmetry breaking (PQ symmetry)

### Pre-inflationary PQ breaking

- Universe = single patch of constant  $\theta$  stretched out by inflation
- Initial axion field value is random 😞
- Inflation dilutes away topological defects 😊

# Axion dark matter – predictions

Axion = pNGB from U(1) symmetry breaking (PQ symmetry)

## Pre-inflationary PQ breaking

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## Post-inflationary PQ breaking

- Universe = huge number of causally disconnected axion field patches
- Axion DM density from realignment = average 😊
- Contribution from top. defects, very difficult to compute 😞 [2007.04990, 2108.05368](#)

- QCD axion mass from chiral perturbation theory<sup>1812.01008</sup>

$$m_a = 5.69(5) \mu\text{eV} \left( \frac{10^{12} \text{ GeV}}{f_a} \right)$$

- Axion-photon coupling depends on UV model through anomaly ratio  $E/N$  and axion-meson mixing<sup>1511.02867</sup>

$$g_{a\gamma\gamma} = \frac{\alpha_{\text{EM}}}{2\pi f_a} \left[ \frac{E}{N} - 1.92(4) \right] \propto m_a$$

- Axion-like particles (ALPs): no connection to QCD = less predictable; however, e.g. mass spectra in string theory<sup>2103.06812</sup>

## Short summary

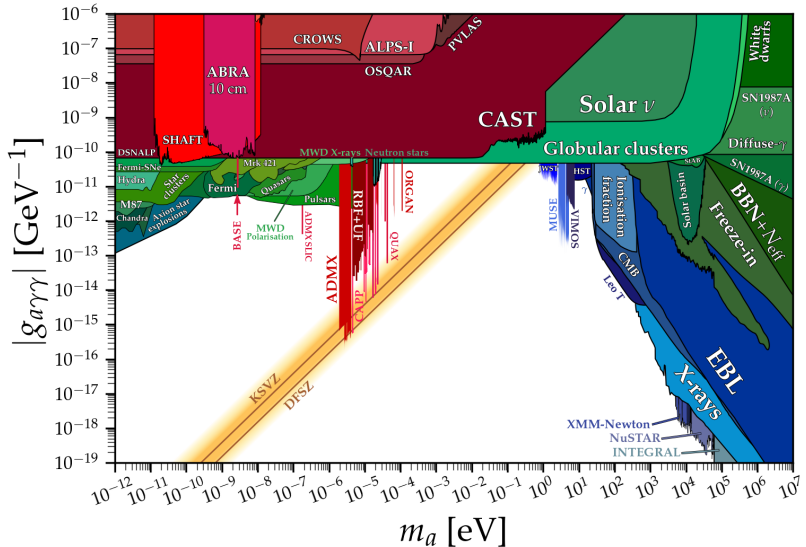
- Axions solve strong  $CP$  problem, explain smallness of  $\theta G\tilde{G}$  term Peccei & Quinn '77 by promoting  $|\theta| \lesssim 10^{-10}$  to a dynamical field
- Unintended bonus: excellent dark matter (DM) candidates!
- Success of axions doesn't depend on PQ scale  $\sim f_a$  😊 But what's the axion's mass? Where to find it? 😞

## Short summary

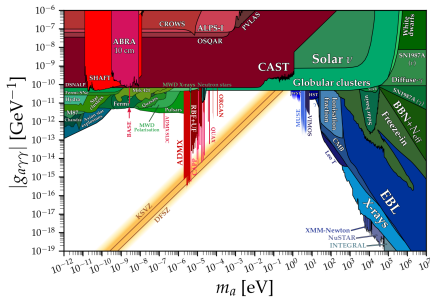
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- Success of axions doesn't depend on PQ scale  $\sim f_a$  😊 But what's the axion's mass? Where to find it? 😞
- String theory: potentially many axion-like particles (ALPs)
- Related ideas: relaxion,<sup>1504.07551</sup> SMASH model,<sup>1610.01639</sup> ALP cogenesis,<sup>2006.04809</sup> ...



# Current limits on the axion-photon coupling



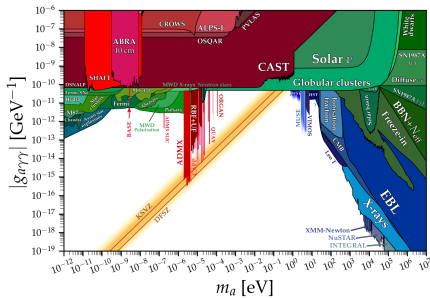
# Global fits for DM ALPs



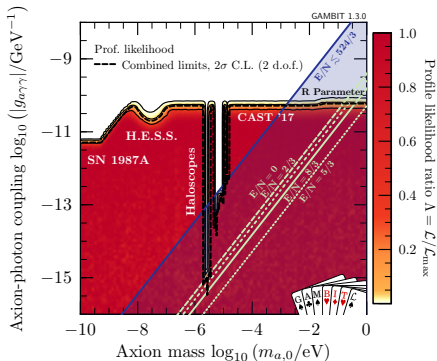
- Consistency of assumptions?
- Overplotted, not combined
- Effects of “hidden parameters”?
- $g_{\gamma\gamma}$  = pheno parameter; no connection to UV model



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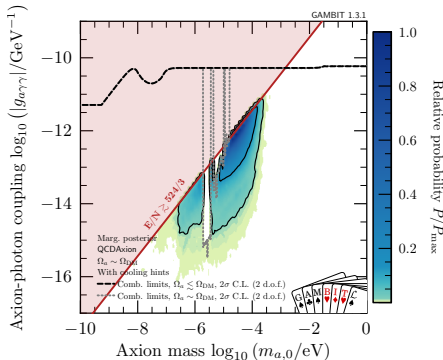
With a global fit, we ensure<sup>1810.07192</sup>

- self-consistent combination and analysis of data
- likelihoods can include all model and nuisance parameters

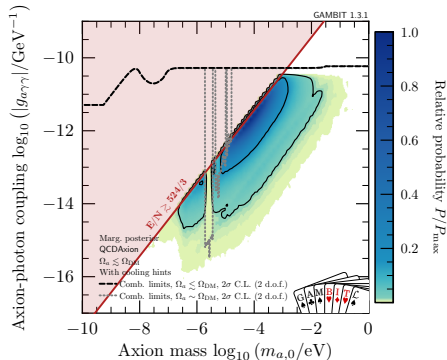
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# Global fits for QCD axions

Where are the most probable, natural QCD axion models in the pre-inflationary PQ breaking scenario? ➔ Bayesian analysis



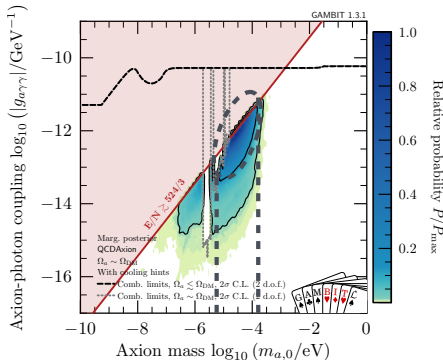
QCD axions = DM



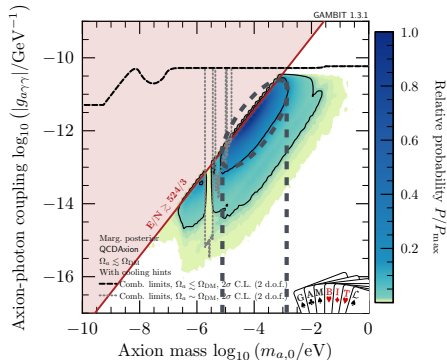
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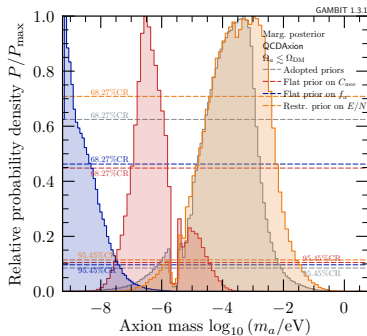
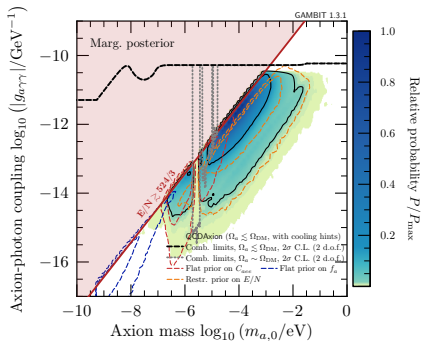
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# Prior dependence of the results

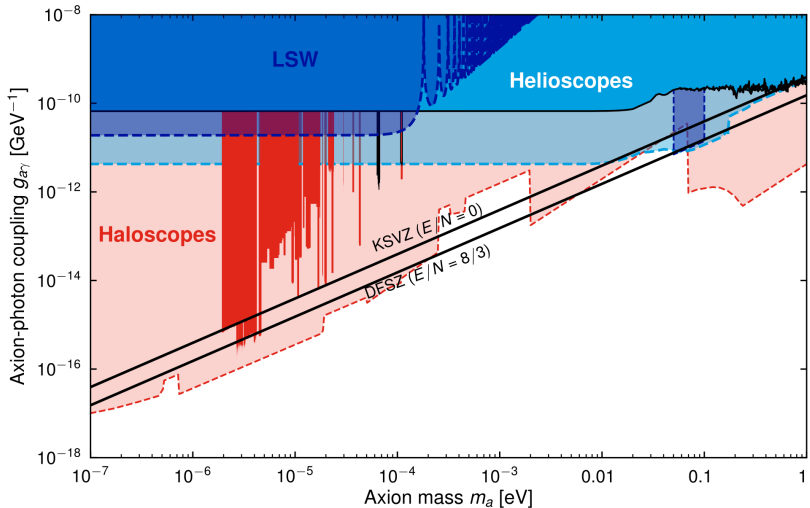
Investigate the prior dependence of the results:



- Uniform vs log uniform priors give very different results
- Are there any “physically motivated” priors?

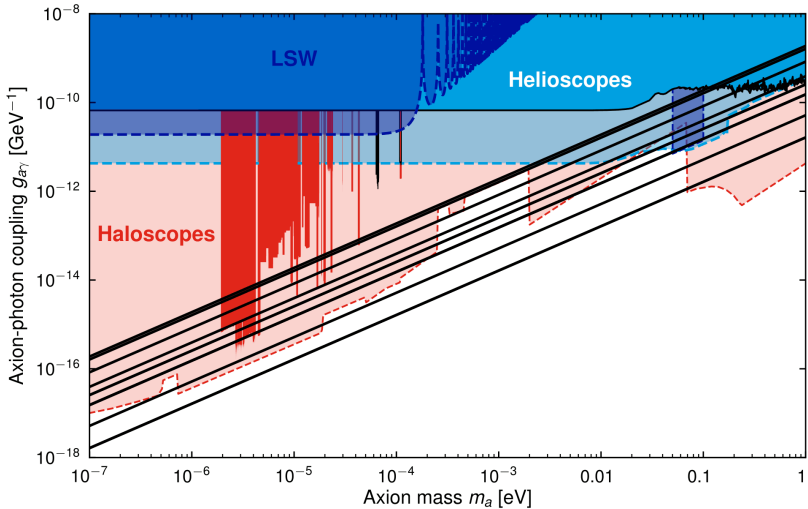
# Aside #1: The genesis of the QCD axion model band

In the beginning there were KSVZ and DFSZ models ...



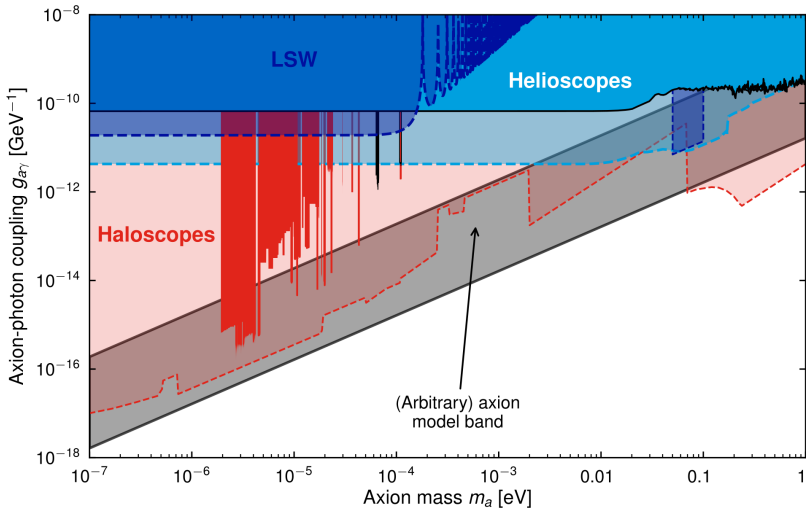
# Aside #1: The genesis of the QCD axion model band

... and then theorists found more models e.g. hep-ph/9506295 ...



# Aside #1: The genesis of the QCD axion model band

... and experimentalists said "let there be a band!" e.g. [hep-ex/0702006](https://arxiv.org/abs/hep-ex/0702006)





- Prior dependence: how to define the “QCD axion band”? Just add more and more models from the vast landscape?<sup>2003.01100</sup>
- Are there infinitely many discrete lines/models? Is the band effectively continuous due to QCD uncertainties?
- Beyond Bayesian analysis and priors: it would just be useful to have a catalogue of models.

Let's start with KSVZ models:

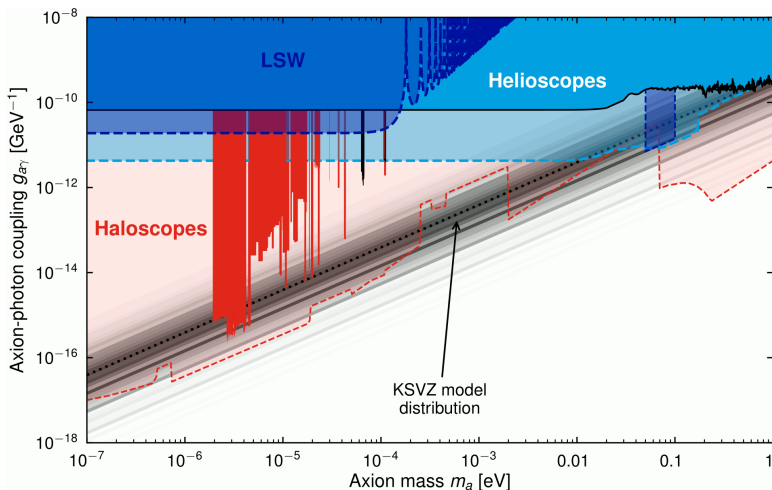
- KSVZ models introduce one new heavy, chiral quark  $Q$ , charged under PQ; charge assignments determine  $E/N$
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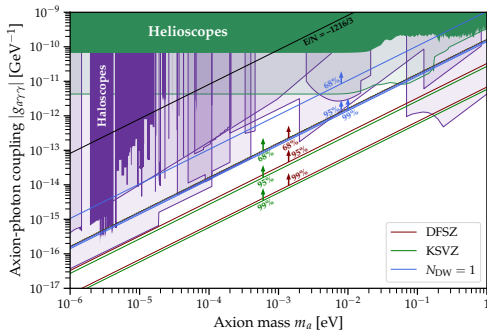
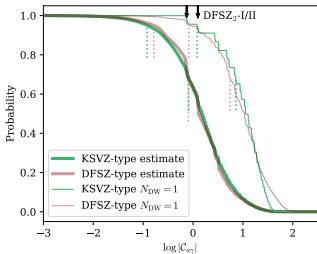
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- Multiple  $Q$ s:  $E/N = (\sum_i E_i) / (\sum_i N_i)$
- Adding too many  $Q$ s will lead to LP below  $m_{\text{Pl}}$ ; gives a *finite number of models*<sup>2107.12378</sup>
- ➔ Creating a (finite) catalogue = combinatorial exercise with selection criteria<sup>1610.07593, 1705.05370</sup>
- N.B.  $N = 0$  possible and the axion does not solve the strong CP problem! *New selection criterion:  $N \neq 0$* <sup>2107.12378</sup>

# The KSVZ model band

Define distribution of “all” KSVZ models (here: equally probable preferred reps)<sup>2107.12378</sup>  $\rightarrow$  theory prior on  $|g_{a\gamma\gamma}| \propto |E/N - 1.92(4)|$

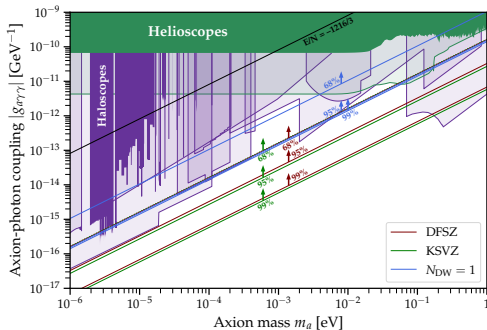
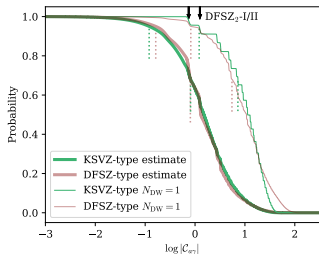


# QCD model band



- Also: DFSZ catalogue available!<sup>2302.04667</sup> Both our<sup>Zenodo</sup> and their<sup>Zenodo</sup> catalogues can be found online
- Discrete  $E/N$  distribution + uncertainties from 1.92(4) term

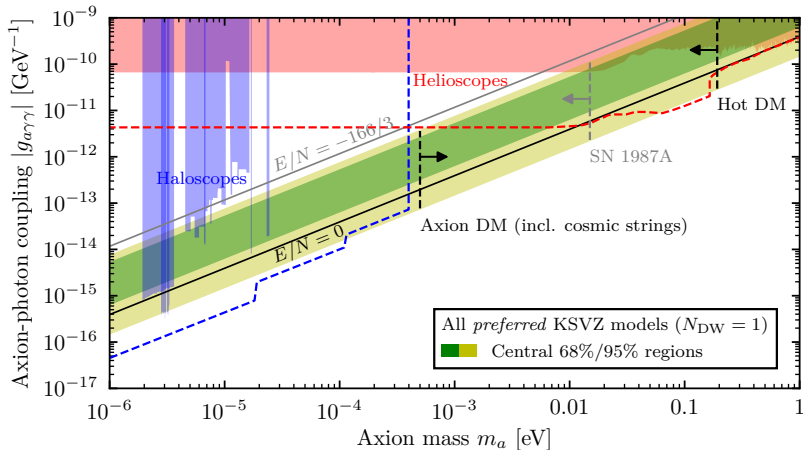
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- ➔ Catalogues can be interpreted as “theory priors” on  $E/N$ , or simply as a database to identify UV models (wait for slide 27)

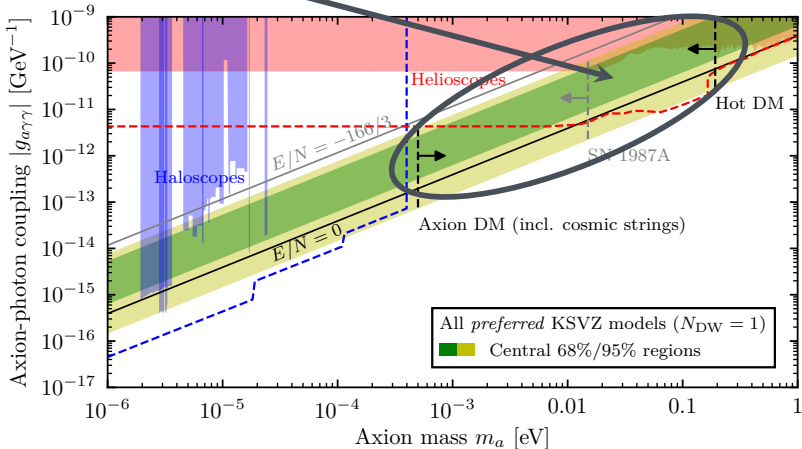
## Example for “boxing in” the axion

Select  $N_{\text{DW}} = 1$  models from KSVZ catalogue (avoids “DW problem” in post-inflationary PQ breaking).



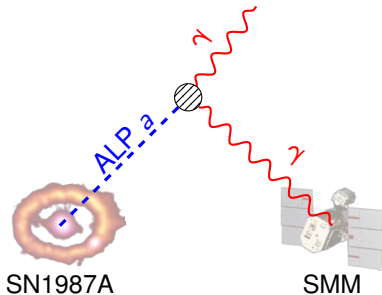
## Example for “boxing in” the axion

Select  $N_{\text{DW}} = 1$  models from KSVZ catalogue (avoids “DW problem” in post-inflationary PQ breaking). For more reliable axion top. defect and thermal production computations: define and probe the  $m_a$ - $g_{a\gamma\gamma}$  regions



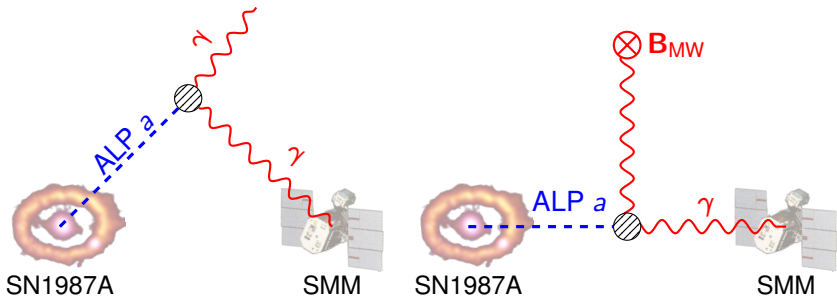


## Aside #2: ALP Constraints from SN1987A



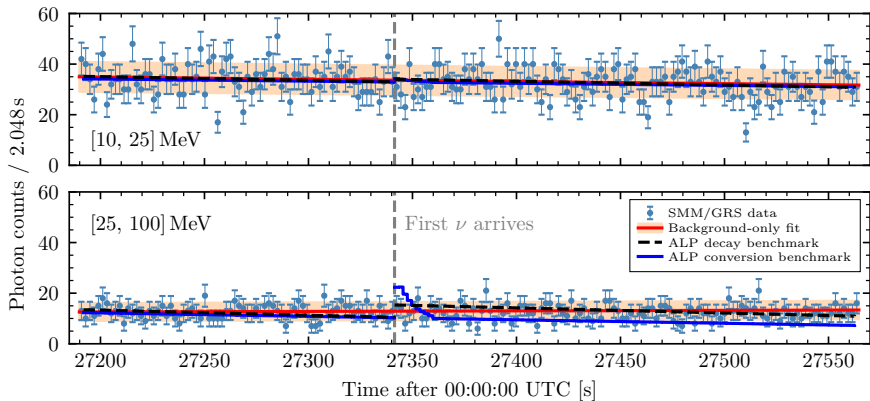
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- We make analytical progress for computations with arbitrary decay lengths; fast code [2212.09764](#)

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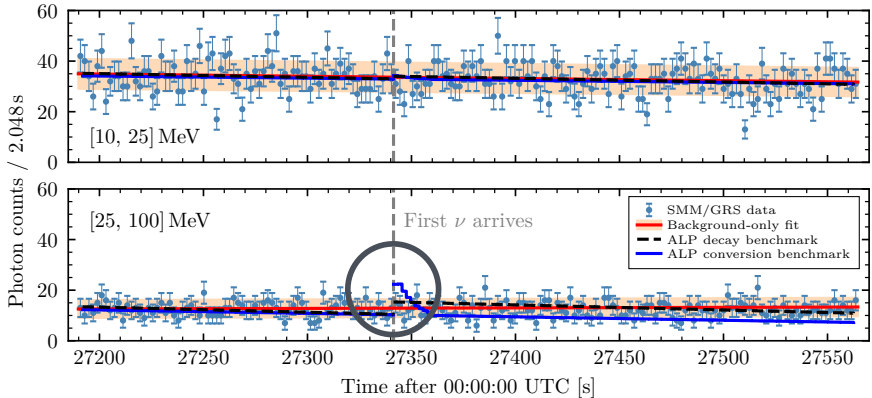
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- We make analytical progress for computations with arbitrary decay lengths; fast code [2212.09764](#)
- Alternatively: Light ALPs can convert to photons in the Galactic  $B$ -field

# SMM data



- Previously: only integrated data was used; no timing info
- Justified for decays: signal is stretched out, approx. const.

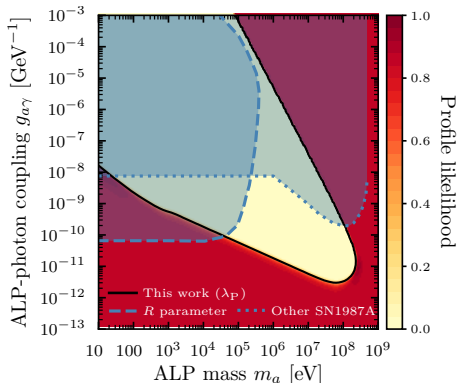
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- Justified for decays: signal is stretched out, approx. const.
- For decays:  $\frac{dN_\gamma}{dt} \propto \frac{dN_a}{dt}$ , time dependence is manifest

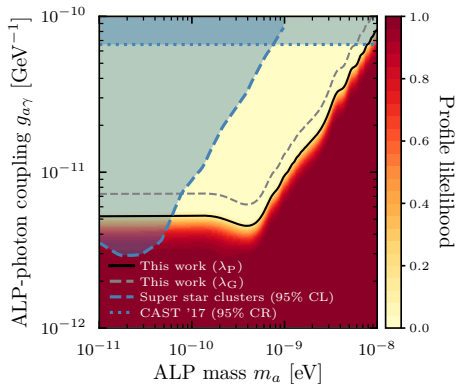
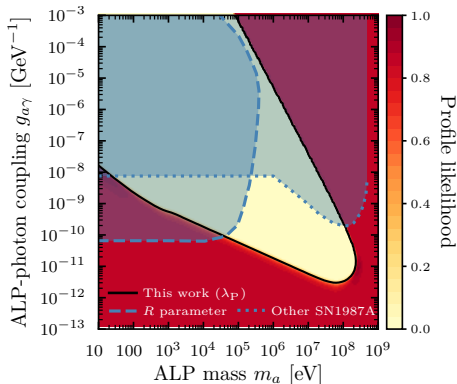
# Updated exclusion limits for SN1987A

- ALP decays: only slight improvement due to additional energy bin, but no significant change (signal  $\approx$  constant)<sup>2212.09764</sup>

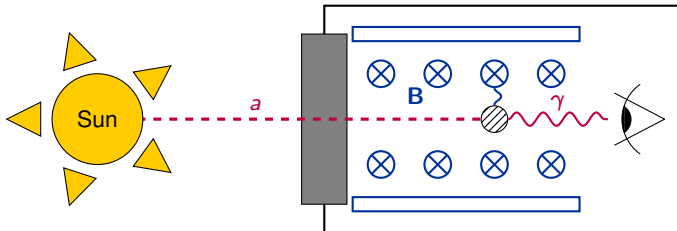


# Updated exclusion limits for SN1987A

- ALP decays: only slight improvement due to additional energy bin, but no significant change (signal  $\approx$  constant)<sup>2212.09764</sup>
- ALP conversions: **factor 1.4 stronger limits**  $\rightarrow$  “global fitting mindset” can help to get more out of the data



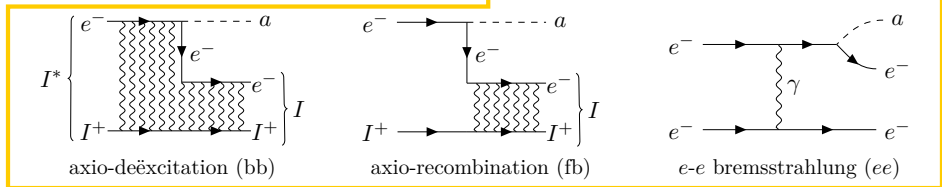
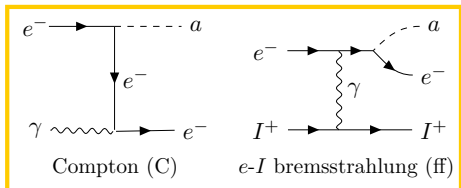
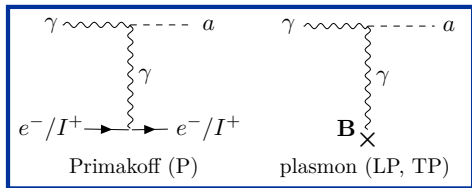
## Helioscopes: detecting solar axions



- $T_{\odot} \sim \text{keV}$ : produce (relativistic) axions in solar plasma
- Axions escape the solar interior  $\approx$  unimpeded
- ➔ Track the Sun across the sky with  $B$ -field + X-ray detector

# Axions production in the Sun

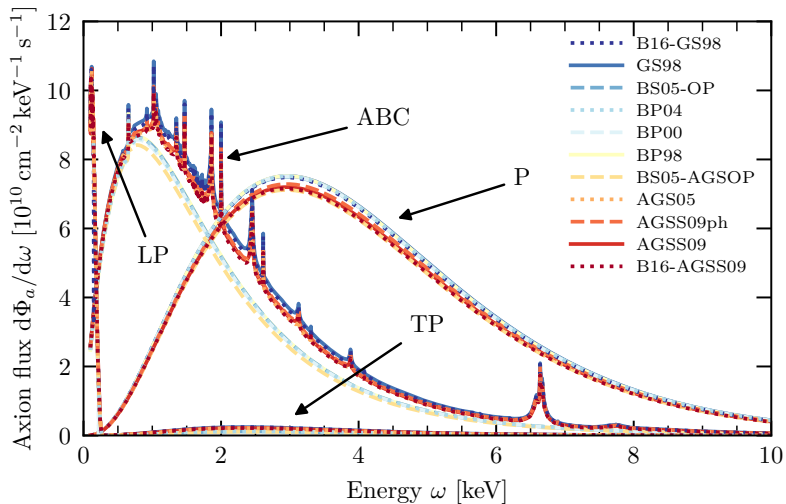
$$\mathcal{L}_{\text{ALP}} = \frac{(\partial_\mu a)^2}{2} - \underbrace{\frac{m_a^2 a^2}{2}}_{m_a \ll T_\odot} - \underbrace{\frac{g_{a\gamma\gamma}}{4} a F\tilde{F} + \frac{g_{aee}}{2m_e} (\partial_\mu a) \bar{e}\gamma^\mu\gamma^5 e}_{[2101.08789]} + \underbrace{\mathcal{L}_{\text{nucl}}}_{[2111.06407]} + \mathcal{L}_{\text{CP}}$$



Production =  $g_{a\gamma\gamma}$ ,  $g_{aee}$ , ..., Raffelt+ '88, ..., Redondo '13, ... detection:  $g_{a\gamma\gamma}$  Sikivie '83 23



# Predictions from solar models

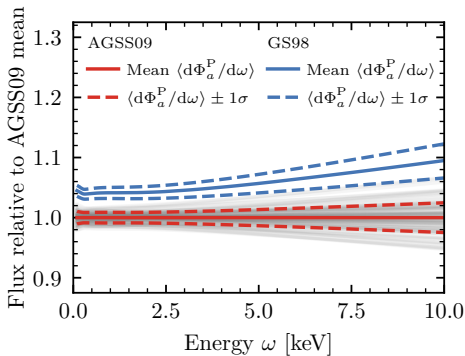


# Solar axion flux uncertainties

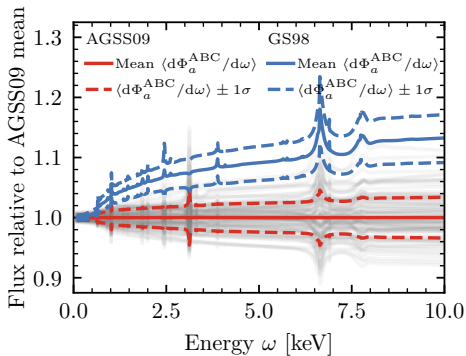
10,000 Monte Carlo sims of low-Z (AGSS09) & high-Z (GS98)

solar models [astro-ph/0511337 + Serenelli update](#) to estimate uncertainties [2101.08789](#)

Primakoff fluxes



ABC fluxes

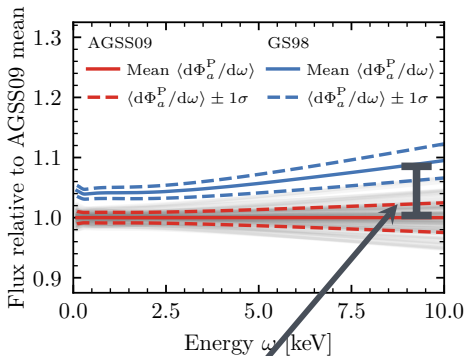


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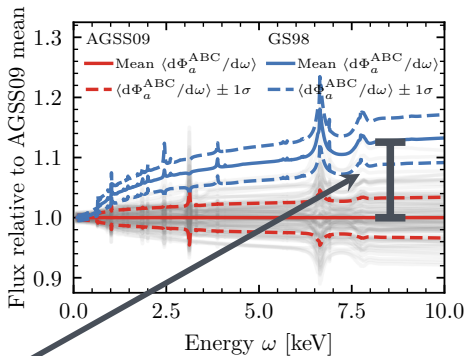
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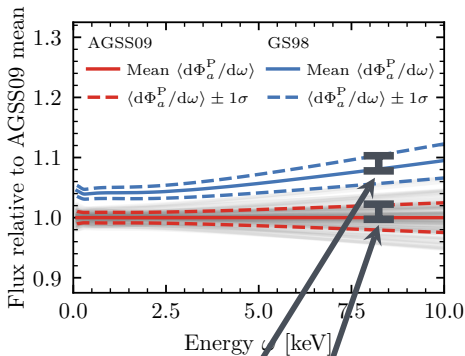
Systematic shift between low-Z and high-Z models (metallicity problem)

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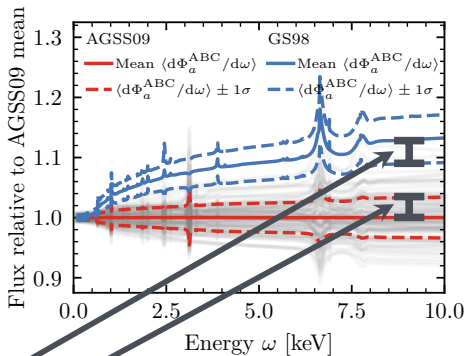
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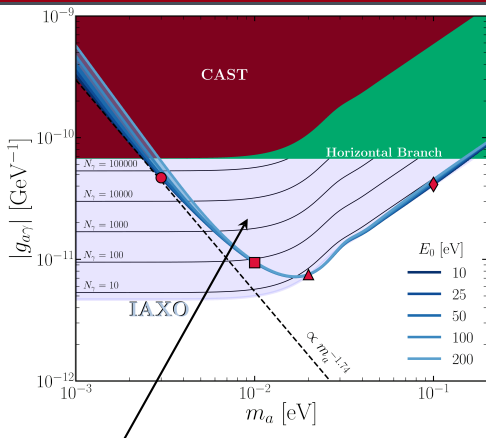


ABC fluxes



Statistical fluctuations; similar for low-Z and high-Z models, smaller than systematics

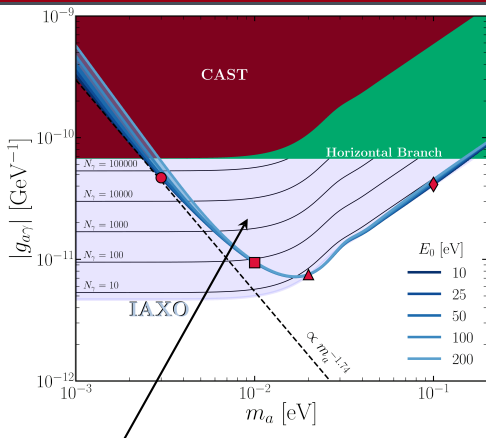
# Discovery potential of IAXO



Parameter regions where IAXO detects  $m_a$  &  $g_{a\gamma\gamma}$  with  $> 3\sigma$  significance, given energy resolution  $E_0$  [1811.09290](#)

- IAXO = helioscope experiment under construction at DESY, Hamburg [1401.3233, 2010.12076](#)
- Can determine  $m_a$  and  $g_{a\gamma\gamma}$  for the region of parameter space on the left

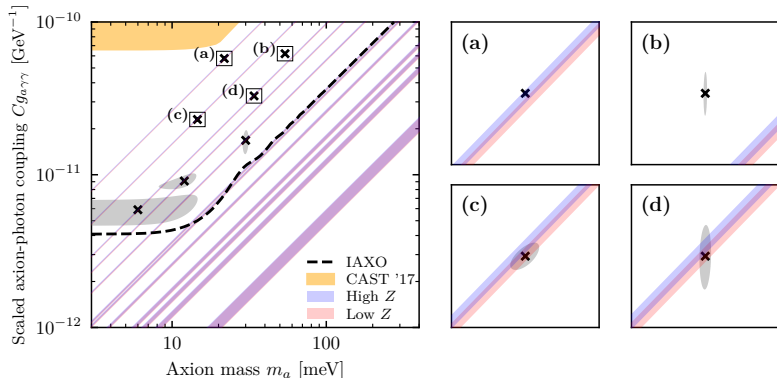
# Discovery potential of IAXO



Parameter regions where IAXO detects  $m_a$  &  $g_{a\gamma\gamma}$  with  $> 3\sigma$  significance, given energy resolution  $E_0$  [1811.09290](#)

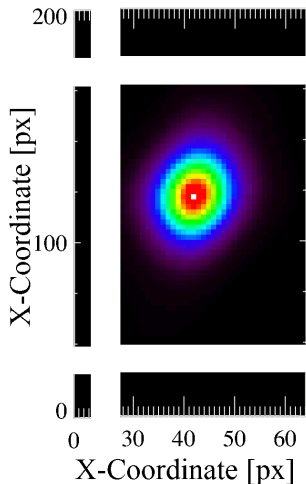
- IAXO = helioscope experiment under construction at DESY, Hamburg [1401.3233](#), [2010.12076](#)
- Can determine  $m_a$  and  $g_{a\gamma\gamma}$  for the region of parameter space on the left
- ➔ Opportunity to discover realistic QCD axion models!
- ➔ Exciting prospect of post-discovery physics

## Other use cases: QCD axion models



- May simultaneously distinguish QCD axion and solar models, [2101.08789](#) hint for solar metallicity problem solution
- Assume Primakoff flux, 15 KSVZ models (pre-catalogue era)
- Can also determine  $g_{aee}$ , [1811.09278](#) metallicities [1908.10878](#)

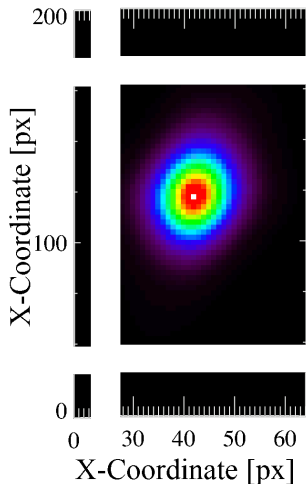
# The solar axion image



- Simulated axion image in CAST helioscope [hep-ex/0702006](https://arxiv.org/abs/hep-ex/0702006)
- $\approx$  spherically symmetric projection thanks to great X-ray optics
- Availability of photon-counting detectors with many pixels

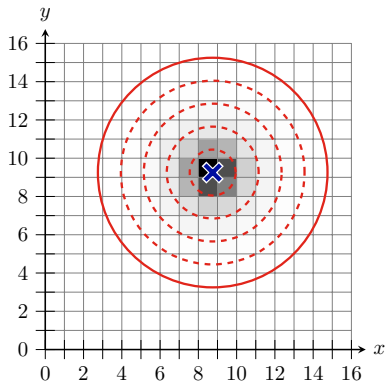


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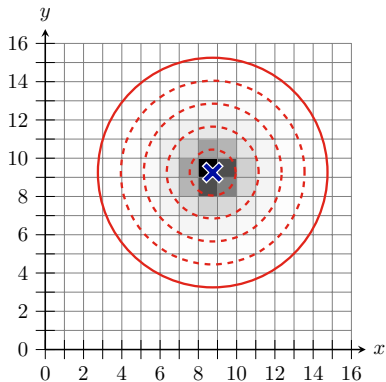
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- $\approx$  spherically symmetric projection thanks to great X-ray optics
- Availability of photon-counting detectors with many pixels
- ➔ Estimate photon counts in rings about the centre of the signal region to obtain radial information

## The solar axion image



- Expected idealised signal in IAXO (actually  $128 \times 128$  pixels, 20 radial, 4 spectral bins)
- Many pixels: photon counts/pixel  $\approx$  equally distributed, integrate flux over radial bins

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- Expected idealised signal in IAXO (actually  $128 \times 128$  pixels, 20 radial, 4 spectral bins)
- Many pixels: photon counts/pixel  $\approx$  equally distributed, integrate flux over radial bins
- Generate 1000 pseudodata sets for IAXO, “invert” solar axion image, fit axion and solar model parameters

## The (simplified) Primakoff production rate

$$\Gamma^{\text{P}}(E_a) = \frac{g_{a\gamma\gamma}^2 \kappa_s^2 T}{32\pi} \left[ \left( 1 + \frac{\kappa_s^2}{4E_a^2} \right) \log \left( 1 + \frac{4E_a^2}{\kappa_s^2} \right) - 1 \right] \frac{2}{e^{E_a/T} - 1}$$

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- Ignores  $e^-$  degeneracy and other corrections (few %)
- ➔ Can break parameter degeneracies with spectral information!

$$\bar{n}_{ij} \propto \int_{\rho_i}^{\rho_{i+1}} d\rho \int_{\rho}^1 dr \frac{r \rho}{\sqrt{r^2 - \rho^2}} \underbrace{\left( \int_{\omega_j}^{\omega_{j+1}} d\omega \frac{\omega^2}{2\pi^2} \Gamma^{\text{P}}(r, \omega) \right)}_{\equiv \bar{\Gamma}_j^{\text{P}}(r)}$$

## A simple reconstruction example

Piecewise-constant interpolation for  $\bar{\Gamma}_j^P$

$$\bar{\Gamma}_j^P(r) = \sum_i \underbrace{\left( \int_{\omega_j}^{\omega_{j+1}} d\omega \frac{\omega^2}{2\pi^2} \Gamma^P(r_i, \omega) \right)}_{\gamma_{i,j}} \Theta(r - r_i) \Theta(r_{i+1} - r)$$

## A simple reconstruction example

Piecewise-constant interpolation for  $\bar{\Gamma}_j^P$  + compute the  $\bar{n}_{i,j}$  integral

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$$\begin{aligned} \bar{n}_{i,j} &\propto \int_{r_i}^{r_{i+1}} d\rho \rho \sum_{k=1}^{n_\rho} \int_\rho^1 dr \frac{r}{\sqrt{r^2 - \rho^2}} \gamma_{k,j} \Theta(r - r_k) \Theta(r_{k+1} - r) \\ &= \frac{1}{3} \left[ \gamma_{i,j} \Delta_{i+1;i}^3 + \sum_{k=i+1}^{n_\rho} \gamma_{k,j} (\Delta_{k+1;i}^3 - \Delta_{k+1;i+1}^3 + \Delta_{k;i+1}^3 - \Delta_{k;i}^3) \right] \end{aligned}$$

with  $\Delta_{\ell,m}^3 \equiv (r_\ell^2 - r_m^2)^{3/2}$

➔ Can compute  $\bar{n}_{i,j}$  analytically!



## A simple reconstruction example

We write this as a matrix equation  $\bar{n}_{i,j} = \sum_{k=1}^{n_p} \mathcal{M}_{ik} \gamma_{k,j}$  with

$$\mathcal{M}_{ik} \propto \begin{cases} \Delta_{i+1;i}^3 & \text{for } i = k, \\ \Delta_{k+1;i}^3 - \Delta_{k+1;i+1}^3 + \Delta_{k;i+1}^3 - \Delta_{k;i}^3 & \text{for } k > i, \\ 0 & \text{otherwise.} \end{cases}$$

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➔ Triangular matrix: set expected = observed counts, invert

$$n_{i,j} = \mathcal{M}_{ii} \gamma_{i,j} + \sum_{k=i+1}^{n_p} \mathcal{M}_{ik} \gamma_{k,j} \Rightarrow \gamma_{i,j} = \frac{1}{\mathcal{M}_{ii}} \left( n_{i,j} - \sum_{k=i+1}^{n_p} \mathcal{M}_{ik} \gamma_{k,j} \right)$$

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➔ Can also propagate errors; use when fitting  $g_{a\gamma\gamma}$ ,  $T_i$  and  $\kappa_i$

$$\sigma_{i,j}^2 \equiv (\Delta \gamma_{i,j})^2 = \frac{1}{\mathcal{M}_{ii}^2} \left[ n_{i,j} + \sum_{k=i+1}^{n_p} \mathcal{M}_{ik}^2 \sigma_{k,j}^2 \right]$$

## Reconstruction in practice

- Matrix only invertible if  $n_{i,j} \neq 0 \Rightarrow$  uneven bin sizes 😞
- More accurate approx. of  $T(r)$  with splines? Ringing 😞

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$$\bar{\Gamma}_j^P(r) = \sum_i \left[ \gamma_{i,j} + \sum_{k=1}^3 c_{k;i,j} (r - r_i)^k \right] \Theta(r - r_i) \Theta(r_{i+1} - r).$$

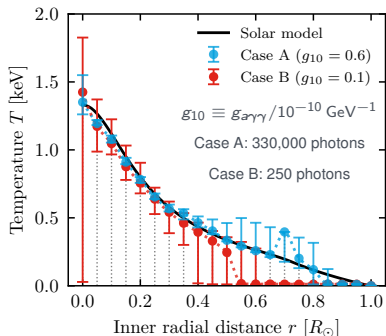
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- Matrix not square, no inversion 😞
- ➔ Direct fitting needed to infer  $g_{a\gamma\gamma}$ ,  $T_i$  and  $\kappa_i$  from the generated pseudodata  $n_{i,j}$ . Optimise:

$$\Delta\chi^2 \equiv -2 \log L(g_{a\gamma\gamma}, \{\kappa_i, T_i\}) = 2 \sum_j \bar{n}_{i,j} - n_{i,j} \log(\bar{n}_{i,j})$$

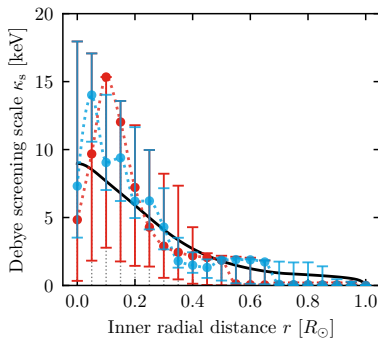
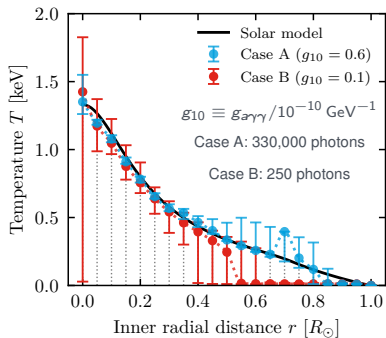
# Temperature reconstruction



We find<sup>2306.00077</sup>

- Accurate  $T(r)$  reconstruction up to  $0.5 R_{\odot}$  ( $0.8 R_{\odot}$ )
- Expected median statistical errors of 10% (16%)

# Temperature reconstruction



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- Accurate  $T(r)$  reconstruction up to  $0.5 R_{\odot}$  ( $0.8 R_{\odot}$ )
- Expected median statistical errors of 10% (16%)
- Difficulties for  $\kappa_s$ : shallow minima, weaker functional dependence, approximation used for  $\Gamma^P$



The upcoming IAXO helioscope can...

... probe more realistic QCD axion models than CAST

... determine mass & couplings<sup>1811.09278, 1811.09290</sup>, *simultaneously distinguish QCD axion and solar models*<sup>2101.08789</sup>

... measure solar metallicities<sup>1908.10878, 2101.08789</sup>

... solar  $B$ -field (profiles),<sup>2005.00078, 2006.12431, 2010.06601</sup>

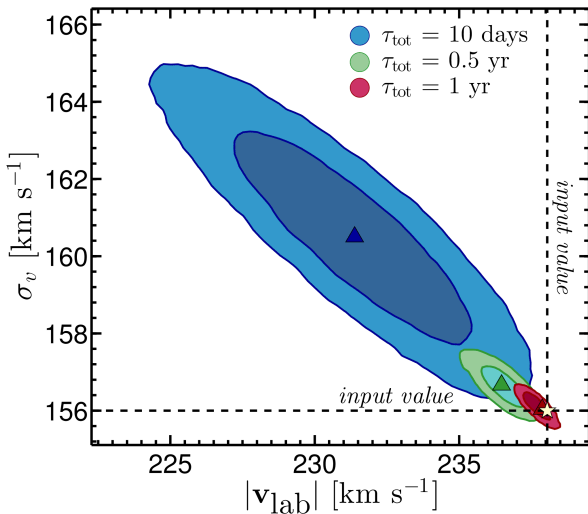
... *measure the solar temperature profile*<sup>2306.00077</sup>

- Axion haloscopes = cavity experiments, tuning the resonant frequency until it matches  $m_a$
- The observed axion power spectrum  $|\mathcal{A}(\omega)|^2$  depends on speed distribution in lab frame  $f_L$ :

$$|\mathcal{A}(\omega)|^2 = 2\pi \frac{\rho_a}{m_a^2} \frac{dv}{d\omega} f_L(v)$$

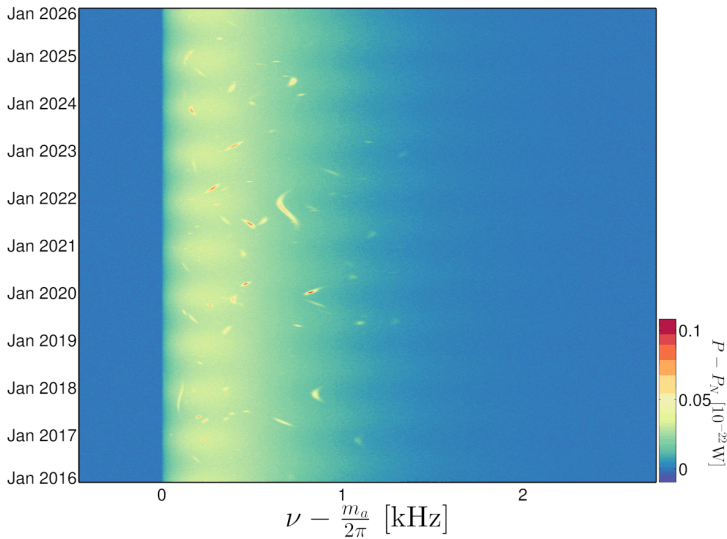
- ➔ Use axions to study local halo properties [1701.03118](#), [1711.10489](#)

## Other post-discovery uses: axion astrometry



Can determine relative halo speed and its dispersion <sup>1701.03118</sup>

## Other post-discovery uses: axion astrometry



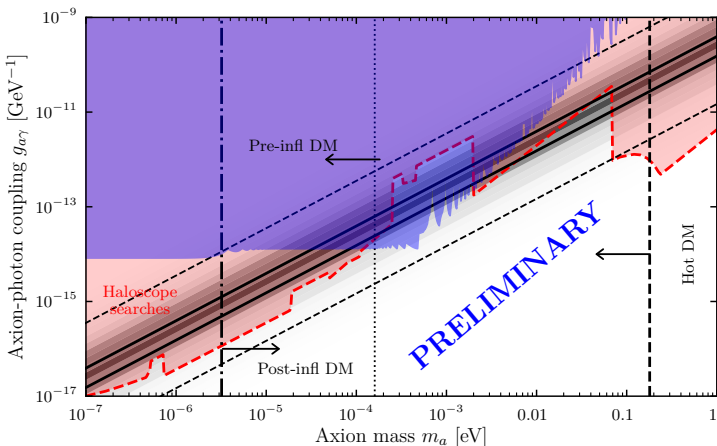
Multi-year obs. can study axion minicluster tidal streams <sup>1701.03118</sup>

- Imagine we find a  $5\sigma$  signal in a haloscope: is it an axions? Is it a QCD axion? What is  $g_{a\gamma\gamma}$ ?
- N.B. we would know  $m_a$  but can only fit  $\rho_{\text{loc}} g_{a\gamma\gamma}^2$
- Can we break the degeneracy? Follow-up experiments needed, but no detailed strategies exist

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- ➔ Use idea for tuning light-shining-through-a-wall experiments with alternating magnet orientations<sup>1009.4875</sup>




# Disentangling axion parameters after a discovery – HyperLSW

*Enter HyperLSW*<sup>†</sup> Total addressable parameter space = union of many individual, tuned magnet arrangements; this only works if you know  $m_a$  since the resonance is narrow!



<sup>†</sup>Working title. Ongoing project w/ J. Jaeckel & G. Lucente

## Summary

- Axions can solve the strong  $CP$  problem, explain DM
- Vast model landscape: value of  $m_a$ ? Where to look?
- DFSZ/KSVZ axion catalogues available now!
- Next-gen helioscopes can discover realistic QCD axion models, determine their mass & couplings
- Solar Primakoff flux predicted at %-level: axions = messengers for solar physics (and beyond)
- Example: accurate, model-independent(!) reconstruction of solar temperature profile  $T(r)$  with axions
- Growing range of open-source software tools for axions: GAMBIT , SolarAxionFlux , snax , ...



## **Bonus Slides**

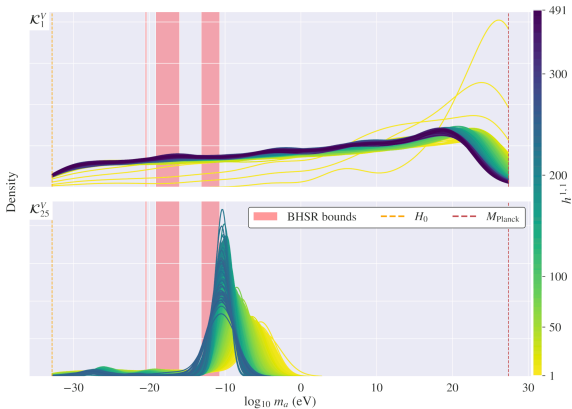
# KSVZ models with one new quark

Repr.	Operator	$E/N$	$N_{\text{DW}}$
$(3, 1, -\frac{1}{3})$	$\bar{Q}_L d_R$	2/3	1
$(3, 1, +\frac{2}{3})$	$\bar{Q}_L u_R$	8/3	1
$(3, 2, +\frac{1}{6})$	$\bar{Q}_R q_L$	5/3	2
$(3, 2, -\frac{5}{6})$	$\bar{Q}_L d_R H^\dagger$	17/3	2
$(3, 2, +\frac{7}{6})$	$\bar{Q}_L u_R H$	29/3	2
$(3, 3, -\frac{1}{3})$	$\bar{Q}_R q_L H^\dagger$	14/3	3
$(3, 3, +\frac{2}{3})$	$\bar{Q}_R q_L H$	20/3	3
$(3, 3, -\frac{4}{3})$	$\bar{Q}_L d_R H^{\dagger 2}$	44/3	3
$(\bar{6}, 1, -\frac{1}{3})$	$\bar{Q}_L \sigma d_R \cdot G$	4/15	5
$(\bar{6}, 1, +\frac{2}{3})$	$\bar{Q}_L \sigma u_R \cdot G$	16/15	5
$(\bar{6}, 2, +\frac{1}{6})$	$\bar{Q}_R \sigma q_L \cdot G$	2/3	10
$(8, 1, -1)$	$\bar{Q}_L \sigma e_R \cdot G$	8/3	6
$(8, 2, -\frac{1}{2})$	$\bar{Q}_R \sigma l_L \cdot G$	4/3	12
$(15, 1, -\frac{1}{3})$	$\bar{Q}_L \sigma d_R \cdot G$	1/6	20
$(15, 1, +\frac{2}{3})$	$\bar{Q}_L \sigma u_R \cdot G$	2/3	20

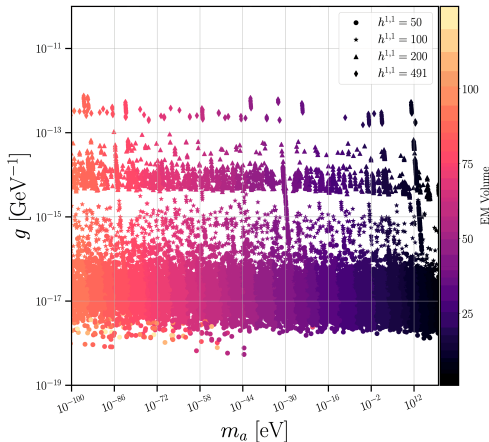
- Define selection criteria for *phenomenologically preferred* models <sup>1610.07593</sup>
- Constraints from lifetimes, DM relic density, Landau poles, ...
- ➔ 15 *preferred* KSVZ-type models with one new exotic quark

# Properties of string theory ALPs

- String theory: many axion-like particles (ALPs) exist<sup>Witten '84, ...</sup>
- How to compute their properties? One approach is to generate random mass matrices etc.<sup>1706.03236, 1909.05257, 2311.13658</sup>
- Recently: explicitly computed mass spectra; can exclude some string theory solutions with BH superradiance?<sup>2103.06812</sup>



# Properties of string theory ALPs



- Even more recently: compute ALP-photon couplings  $g_{a\gamma\gamma}$ , so we can do more phenomenology!<sup>2309.13145</sup>

➔ Q: how do deal with the complexity of multi-ALP theories?

## Multi-ALP systems from string theory

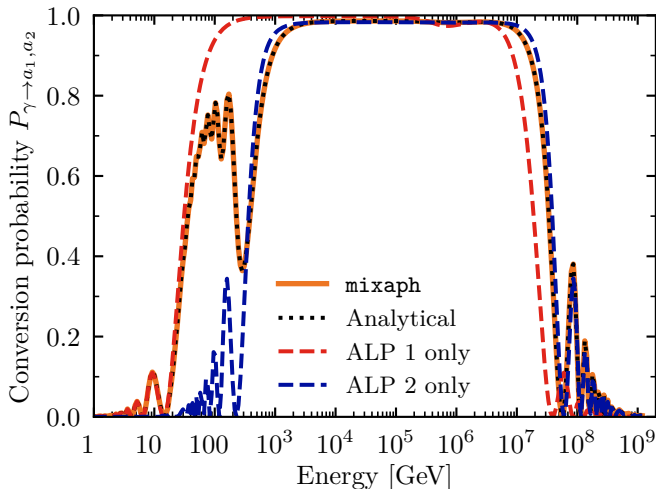
Considering 100s of ALPs (from string theory) is tricky:

- Mass oscillations over longer distances – even without explicit ALP-ALP interactions<sup>2107.12813</sup>
- $\mathcal{H}$  for ALP-photon system  $\approx$  sparse, but grows as  $\propto N_{\text{ALP}}^2$ !
- ➔ Numerical approach needed!

We can make our life easier:

- Can sum up (effectively) massless states and ignore “heavy” states<sup>1909.05257, 2107.12813</sup>
- Relevant length and energy scales will depend on the ALP search<sup>2311.13658</sup>
- ➔ Still need a code to solve a system of  $\mathcal{O}(10)$  ALPs

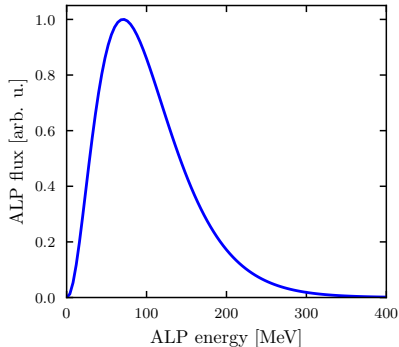
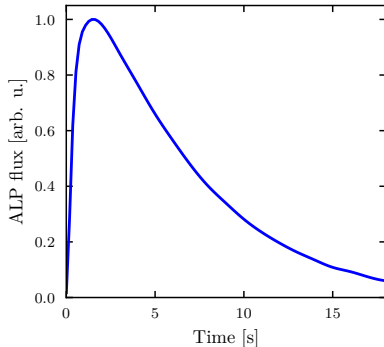
## Mixaph (WIP)



Enter mixaph! Upcoming software code to compute predictions of multi-ALP systems; esp. relevant for astrophysical constraints


# ALP spectrum from SN1987A

- Use the ALP spectrum computed in previous work<sup>1410.3747</sup>
- Rescale cross section to approximate massive case<sup>1702.02964</sup>
- Ignore photon coalescence<sup>2008.04918, 2107.12393</sup>



## Observational constraints on the origin of the elements

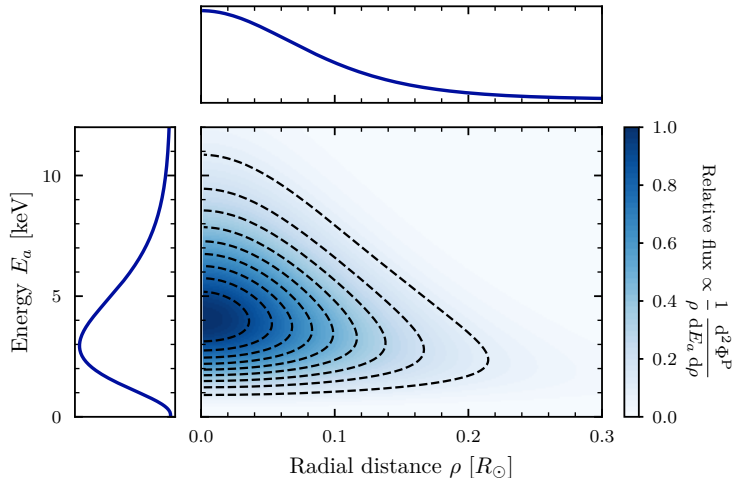
### IV. Standard composition of the Sun

Ekaterina Magg<sup>1</sup>, Maria Bergemann<sup>1,5</sup>, Aldo Serenelli<sup>2,3,1</sup>, Manuel Bautista<sup>4</sup>, Bertrand Plez<sup>7</sup>, Ulrike Heiter<sup>6</sup>, Jeffrey M. Gerber<sup>1</sup>, Hans-Günter Ludwig<sup>8</sup>, Sarbani Basu<sup>9</sup>, Jason W. Ferguson<sup>10</sup>, Helena Carvajal Gallego<sup>11</sup>, Sébastien Gamrath<sup>11</sup>, Patrick Palmeri<sup>11</sup>, and Pascal Quinet<sup>11,12</sup>

- New composition: MB22<sup>2203.02255</sup> (models available now<sup>Zenodo</sup>)
- Claims to reproduce sound velocity profile  $c(r)$  with both photospheric and meteoritic abundances? (However: potential issues?<sup>2308.13368</sup>)
- ➔ Benefits of our open-source code: re-compute all fluxes for models based on new compositions once available

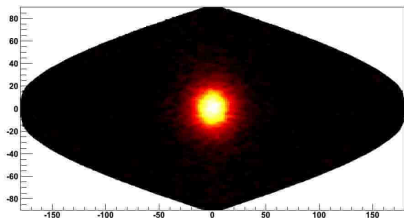


## Primakoff flux on the solar disc



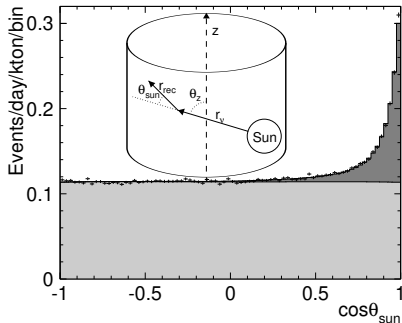
- Primakoff process dominant for KSVZ
- 50% (99%) of P flux contained within  $0.15 R_\odot$  ( $0.5 R_\odot$ )
- Few % stat. and sys. errors

# Can we reconstruct solar $T(r)$ with $\nu$ s?



Super-K Collaboration 1998–2018

Solar  $\nu$  image with more than  $10^5$  events!



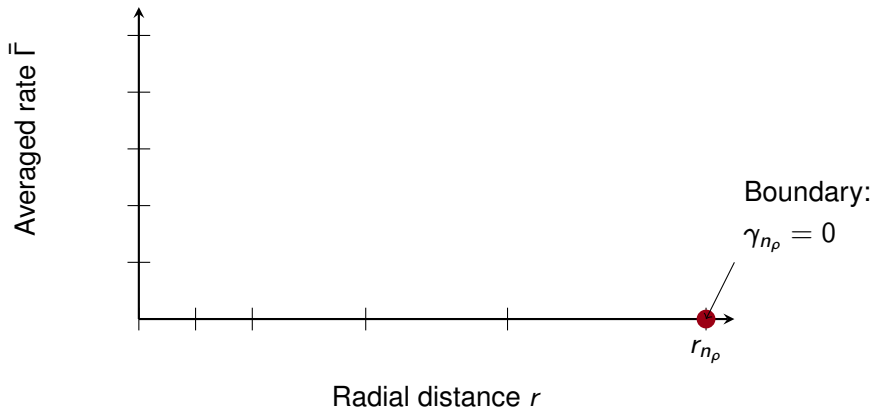
1606.07538

Sadly: angular res.  $\sim 40^\circ$  vs the Sun's apparent size of  $\sim 0.5^\circ$ ,  $e^-$  recoil and  $\nu$  path not aligned

➡ Helioscope X-ray optics offer superior spatial resolution

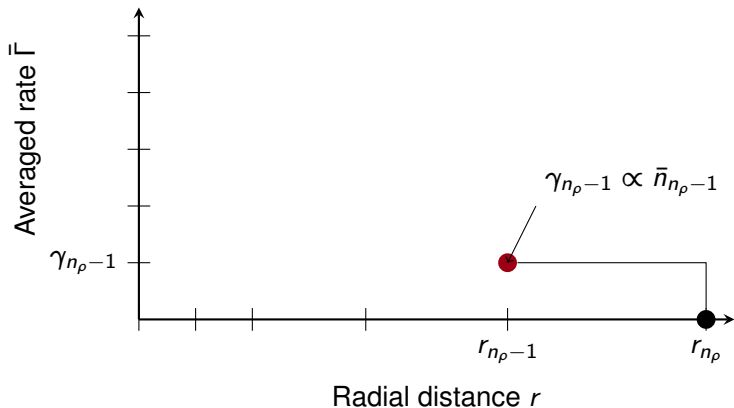
# Reconstruction

For the  $j$ th energy bin, the reconstruction works as follows:



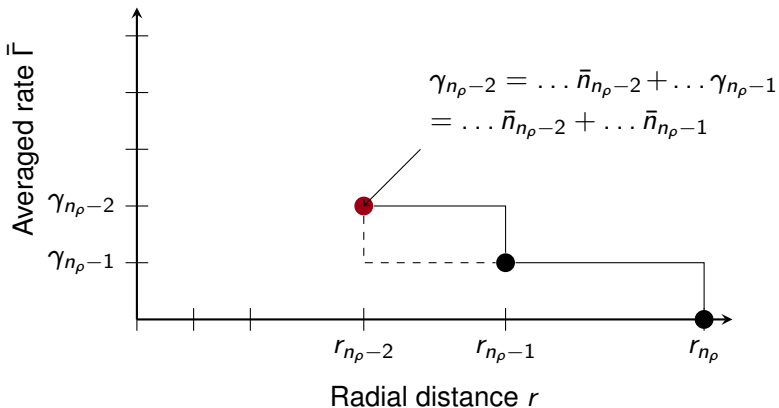
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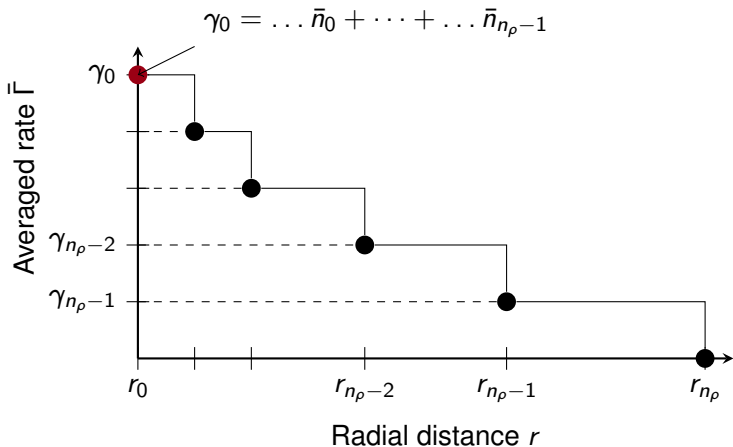
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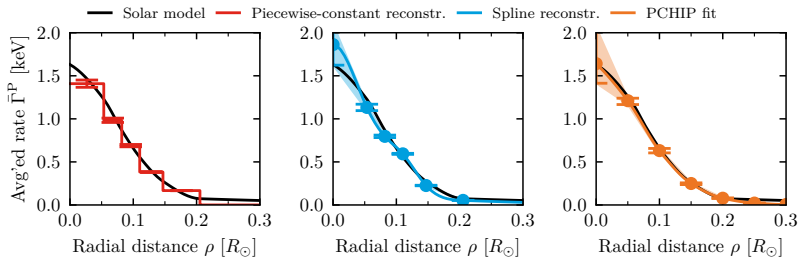


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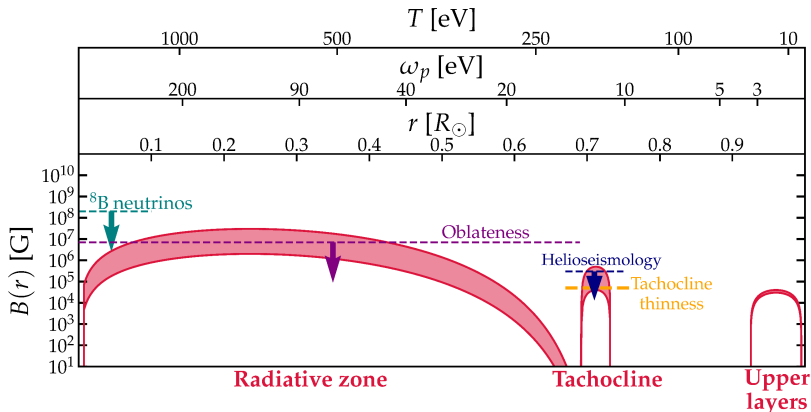
For the  $j$ th energy bin, the reconstruction works as follows:



# Different reconstruction techniques for $T(r)$



# Axions as solar magnetometers



- Axions are produced in macroscopic solar  $B$  fields through plasmon interactions [2005.00078](#), [2006.10415](#), [2010.06601](#)
- Mostly resonant phenomenon: relates  $r \leftrightarrow \omega_{\text{pl}} \leftrightarrow E_a$
- ➔ Can map  $B(r)$  [2006.10415](#) — impossible w/o axions!?